



# Design of photonic microcavities in hexagonal boron nitride

Sejeong Kim<sup>\*</sup>, Milos Toth and Igor Aharonovich

## Letter

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Address:  
School of Mathematical and Physical Sciences, University of  
Technology Sydney, Ultimo, NSW 2007, Australia

Email:  
Sejeong Kim<sup>\*</sup> - Sejeong.Kim-1@uts.edu.au

\* Corresponding author

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## Abstract

We propose and design photonic crystal cavities (PCCs) in hexagonal boron nitride (hBN) for diverse photonic and quantum applications. Two dimensional (2D) hBN flakes contain quantum emitters which are ultra-bright and photostable at room temperature. To achieve optimal coupling of these emitters to optical resonators, fabrication of cavities from hBN is therefore required to maximize the overlap between cavity optical modes and the emitters. Here, we design 2D and 1D PCCs using anisotropic indices of hBN. The influence of underlying substrates and material absorption are investigated, and spontaneous emission rate enhancements are calculated. Our results are promising for future quantum photonic experiments with hBN.

## Introduction

Hexagonal boron nitride (hBN) has recently emerged as an interesting platform for nanophotonics. This is mainly due to its promising hyperbolic properties [1,2] as well as the ability to host a range of single photon emitters (SPEs) that are of great interest for a myriad of nanophotonics and quantum photonic applications [3–10]. However, to further study light matter interactions based on the hBN SPEs, and to realize integrated nanophotonics systems, coupling of the emitters to optical cavities is essential [11–15].

Typically, SPEs can be coupled to optical cavities using two general approaches. The cavity is either made from the material

that hosts the emitter (monolithic approach), or the emitter is coupled to a cavity made from a foreign material – also known as a hybrid approach [16]. The former process is preferred when attempting to maximize the field overlap between the emitters and the cavity modes, and is often employed when using materials that are amenable to scalable nanofabrication protocols, such as gallium arsenide or silicon [17,18], and more recently diamond and silicon carbide [19,20]. The hybrid approach is easier from the fabrication point of view but is inherently limited by the fact that the electric field maxima of optical modes are situated within the cavities, and optimal coupling therefore remains a challenge.

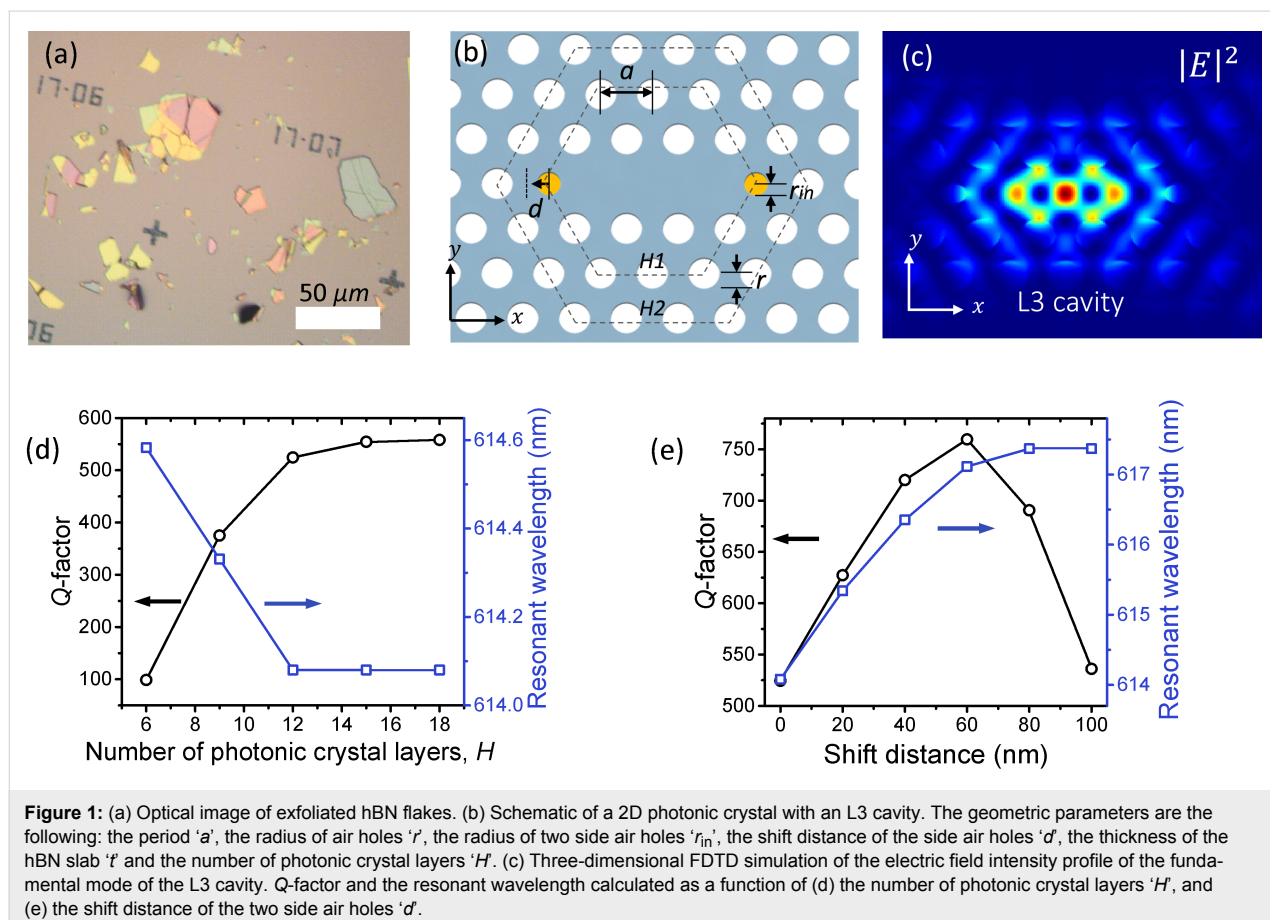
The optical properties of hBN make it an attractive candidate for a monolithic cavity system. In particular, hBN has a wide bandgap of  $\approx 6$  eV which makes it transparent in the visible spectral range that contains the zero phonon lines (ZPLs) of a range of ultra-bright emitters [21]. Furthermore, hBN has properties which are desirable for micro-resonators such as a high chemical stability and an excellent thermal conductivity [22,23].

In this work, we propose to use hBN for the fabrication of photonic crystal cavities (PCCs). We design two dimensional (2D) PCCs and show that they have high quality-factor ( $Q$ -factor) resonances in the visible spectral range, which overlap with the ZPLs of SPEs in hBN [24]. We further optimize the structures and model 1D nanobeam photonic crystals that exhibit a  $Q$ -factor in excess of  $\approx 20,000$ . In the light of recent progress in direct-write etching of hBN [25], our results are promising for realization of high  $Q$  cavities and monolithic coupled systems made from this material.

## Results and Discussion

We begin with a 2D photonic crystal that contains a line defect cavity. The L3 cavity has been widely investigated because it

was the first to exceed an experimental  $Q$ -factor of 10,000 [26,27]. In this study, we used a commercial finite-difference time-domain (FDTD) software package (Lumerical Inc.). The 3D FDTD simulation domain for 2D (1D) photonic crystal was  $7 \mu\text{m} \times 7 \mu\text{m} \times 1.2 \mu\text{m}$  ( $11 \mu\text{m} \times 2 \mu\text{m} \times 2 \mu\text{m}$ ) which is discretised using uniform spatial and temporal grids of 15 nm and 0.03 fs. Birefringence of hBN is accounted for in our study by including both ordinary ( $n_x = n_y = 1.72$ ) and extraordinary ( $n = 1.84$ ) indices in the 3D FDTD method models. Figure 1a shows the shape and the size of typical hBN flakes prepared by scotch-tape exfoliation [9]. The lateral flake size varies from a few micrometers to few tens of micrometers while thickness varies from few tens of nanometers to a few micrometers, which is sufficient for the fabrication of practical photonic crystal cavities. The parameters used to define an L3 cavity are shown in Figure 1b. The cavity consists of a free-standing slab with a triangular photonic lattice with periodicity ' $a$ '. The air hole radius in the mirror region and the radius of two side air holes are fixed at  $0.33a$  and  $0.22a$ , respectively. By tuning the periodicity ' $a$ ', one can tune the resonant wavelength whilst preserving the  $Q$ -factor. In this study, we used  $a = 270$  nm and  $t = 280$  nm to place the fundamental mode within the typical emission range of SPEs in hBN (550–700 nm). Two air holes at

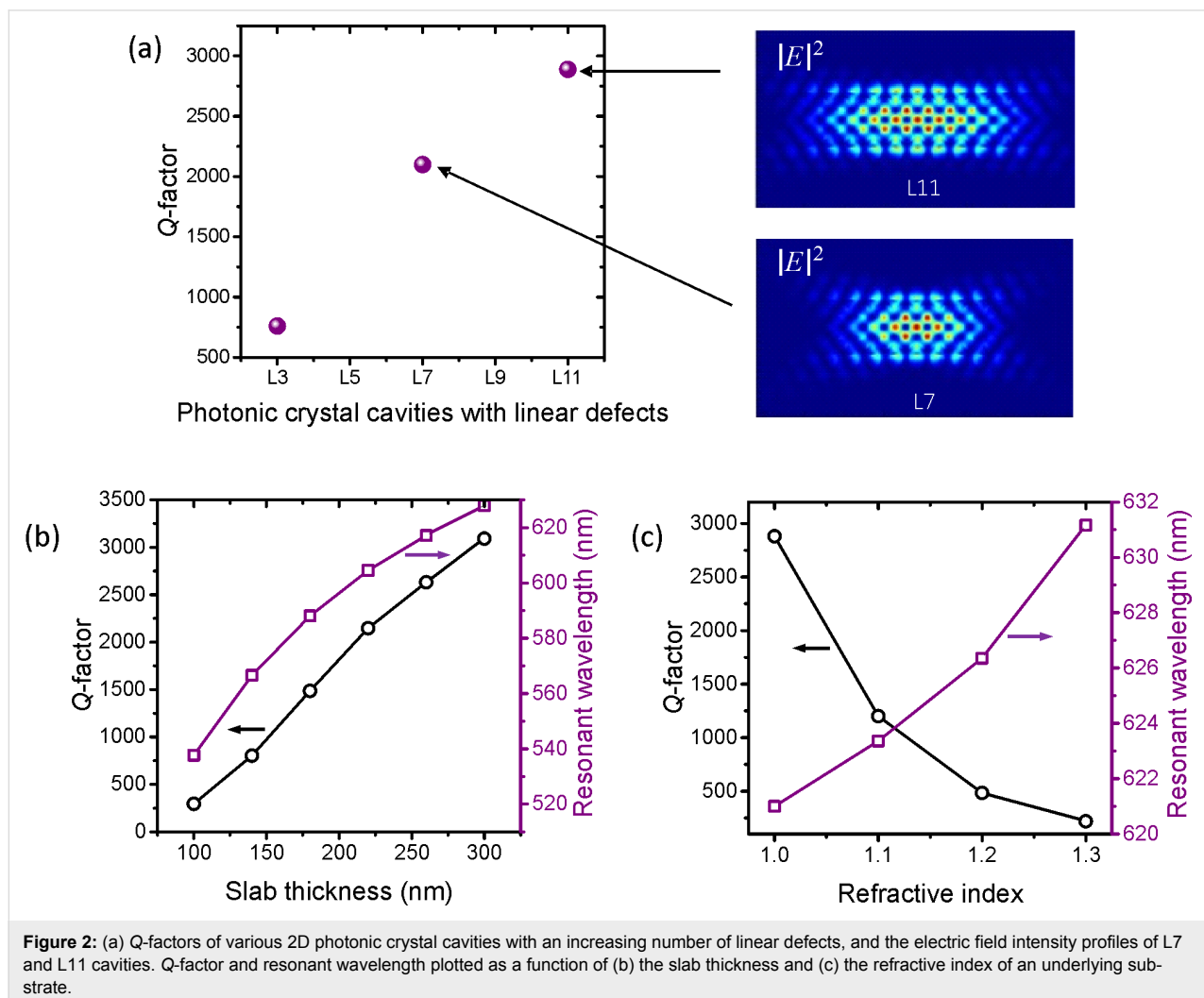


the side of the cavity (yellow circles in Figure 1b) are decreased in size to reduce radiation losses [28]. For the purpose of design simplicity, we adopt tuning of the side air holes as the only means used to increase the  $Q$ -factor. Lastly, the number of photonic crystal layers comprising the cavity is denoted by ' $H$ ' which act as a photonic mirror and the thickness of the hBN slab by ' $t$ '. The electric field intensity pattern of the L3 cavity calculated using 3D FDTD simulation is shown in Figure 1c. This is the lowest energy mode in the L3 cavity which is the most widely studied of high- $Q$  2D photonic cavities.

To optimize the design, we start by increasing the number of photonic crystal layers  $H$ , in order to increase the photonic mirror strength and to reduce the in-plane loss of the cavity. As shown in Figure 1d, the  $Q$ -factor of the mode starts to saturate at  $H \approx 12$  because the  $Q$ -factor is limited not only by the in-plane component, but also by radiation loss. Considering both the  $Q$ -factor and scaling of the simulation time with domain size, we fixed  $H$  at 12 for subsequent modelling.

Radiation losses can be reduced by optimizing the side air hole positions, as is shown in Figure 1e, which reveals that the cavity  $Q$ -factor is greatest at a shift distance of 60 nm. Note that the mode resonance red shifts with increasing side hole separation due to an effective increase in cavity length.

Photonic crystal cavities with line defects can be described as Fabry–Pérot resonators [29]. Hence, the  $Q$ -factor can be enhanced by increasing the cavity length as for a Fabry–Pérot resonator [30], which can be tuned by varying number of missing air holes. Figure 2a shows the  $Q$ -factors of L3, L7, and L11 cavities with a slab thickness of 280 nm. The electric field intensity profiles of L7 and L11 are also shown in the figure (and that of the L3 cavity is shown in Figure 1c). The  $Q$ -factor of the fundamental mode increases with effective cavity length. For the L11 cavity, we additionally calculated the effect of the thickness of the 2D photonic crystal slab as is shown in Figure 2b. The simulated slab thickness is varied up to 300 nm since that is a realistic thickness of typical hBN flakes prepared

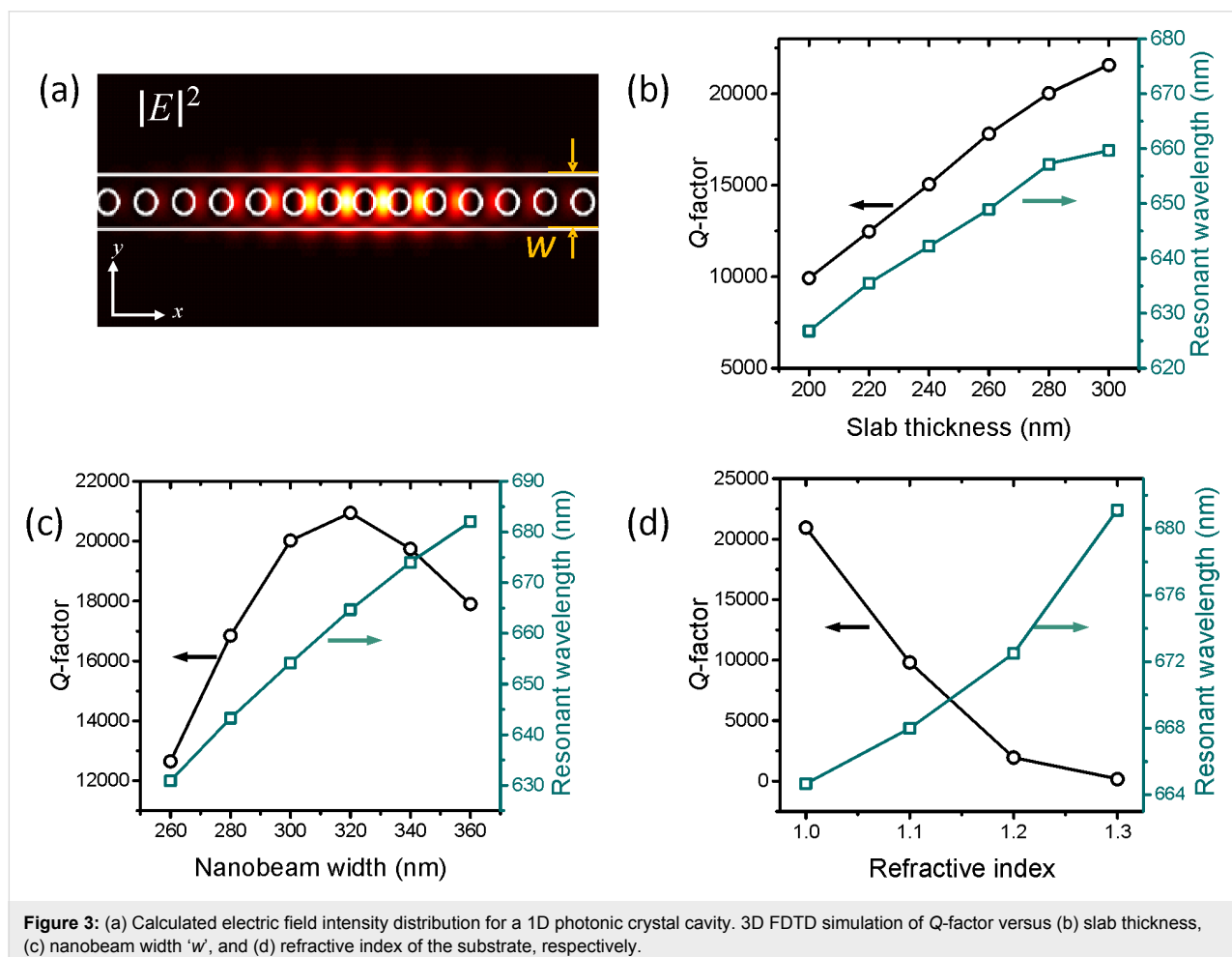


by the scotch tape exfoliation method. Thicker slabs exhibit stronger light confinement, which results in higher  $Q$ -factors.

We also modelled the effect of the refractive index of an underlying substrate, as is shown in Figure 2c. Because the refractive index of hBN is relatively low compared to that of typical semiconductors, the increase in substrate index greatly degrades the  $Q$ -factor of the L11 cavity (and a similar effect is expected for the other cavities as well). This is, however, not a significant problem as the transfer of hBN flakes onto holey substrates is a straightforward process. Furthermore, use of aerogel material that is currently commercially available is another option to achieve low index substrates. We note that the overall  $Q$ -factor is lower than in typical semiconductor 2D photonic crystals due to the relatively low refractive index of hBN.

Next we investigate one dimensional nanobeam photonic crystal structures. 1D PCCs are advantageous in that they can easily have full photonic bandgap between the first and the second lowest photonic energy bands even when the effective index contrast is low [31]. The combination of a high  $Q$ -factor

and a low refractive index enables a broad range of applications such as flexible photonic crystal devices and high figure of merit sensors [32]. Figure 3a shows the electric field intensity profile of the fundamental mode. The cavity is designed by modulating the periodicity whilst fixing the air hole radius [33]. The structure consists of a total of 31 air holes, 15 of which are modulated to create a cavity in the center, and the remaining 8 on each end act as photonic mirrors. In Figure 3b, we set the periodicity in the mirror region to 260 nm while the air hole radius and the nanobeam width are fixed at 70 nm and 300 nm, respectively. By increasing the thickness of the slab from 200 nm to 300 nm, the  $Q$ -factor increases, as is seen in Figure 3b, and  $Q$ -factors in excess of 20,000 can be realized. These values are achieved even without optimization of the remaining structural parameters, and are over an order of magnitude greater than the maximum  $Q$ -factor of a low-index 2D cavity. Next, we fix the thickness at 280 nm and tune the structural parameters to further increase the  $Q$ -factor of the mode. Figure 3c shows the  $Q$ -factor plotted as a function of the nanobeam width ‘ $w$ ’, showing that the  $Q$ -factor has a maximum at a width of 320 nm [34]. Introducing a substrate to the

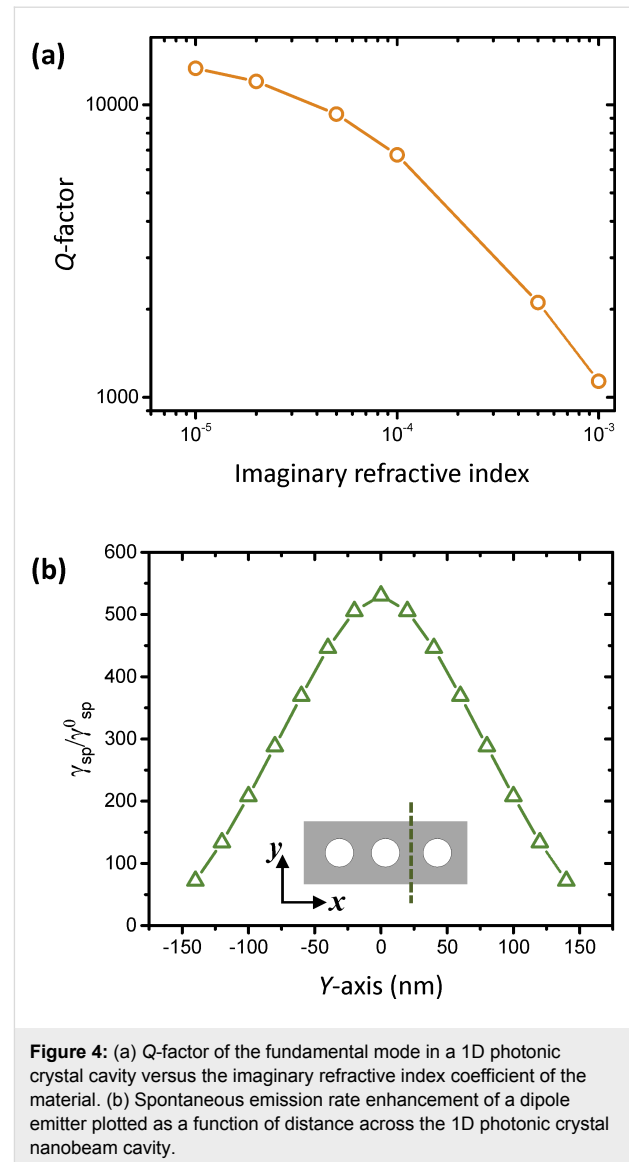


**Figure 3:** (a) Calculated electric field intensity distribution for a 1D photonic crystal cavity. 3D FDTD simulation of  $Q$ -factor versus (b) slab thickness, (c) nanobeam width ‘ $w$ ’, and (d) refractive index of the substrate, respectively.

nanobeam degrades the  $Q$ -factor as is shown in Figure 3d. A free-standing structure is preferred, as in the case of the 2D photonic crystals presented earlier.

We also consider the absorption losses in the cavity. Figure 4a shows the effect of the absorption by the cavity material. Many FDTD modelling studies of semiconductor photonic crystals include only the real part of the complex refractive index as a structural input which assumes that the material is transparent in the simulated wavelength regime. However, in the case of practical situations in which the cavity material exhibits finite absorption, the imaginary part should also be accounted for. The  $Q$ -factor of the fundamental mode in a free-standing nanobeam with a beam-width ( $w$ ) of 320 nm and a slab thickness ( $t$ ) of 280 nm is calculated as a function of the imaginary refractive index. An increase in imaginary refractive index (i.e., an increase in material absorption) causes the  $Q$ -factor to decrease significantly. Note that realistic  $Q$ -factor is determined by  $1/Q = 1/Q_{\text{ideal}} + 1/Q_{\text{abs}}$  ( $Q_{\text{ideal}}$ :  $Q$ -factor with lossless material,  $Q_{\text{abs}}$ :  $Q$ -factor with absorption losses) [35], which indicates the material loss restricts the maximum  $Q$ -factor that can be obtained through experiment. Therefore, including imaginary refractive index for  $Q$ -factor calculation provides practical  $Q$ -factors.

Finally, we discuss the expected Purcell enhancement due to coupling of emitters to cavity modes. Figure 4b shows the spontaneous emission rate ( $\gamma_{\text{sp}}$ ) of a quantum emitter in the nanobeam cavity relative to an emitter in free space ( $\gamma_{\text{sp}}^0$ ). The Purcell enhancement ( $\gamma_{\text{sp}}/\gamma_{\text{sp}}^0$ ) was found at any given spatial position by calculating the radiated power enhancement of the quantum emitter relative to that in free space. Because the fundamental mode of the nanobeam has a maximum in the high index region, the  $Q$ -factor is calculated across the dashed line ( $y$ -axis) shown in the inset (i.e., along the width of the 1D PCC). The dipole emitter in the simulation is  $y$ -polarized to match the polarization of the optical mode. The Purcell enhancement has a maximum in the center of the nanobeam where the electric field intensity is the greatest. The expected Purcell enhancement is greater than 100 over a range of more than 200 nm along the  $y$ -axis, as is seen in Figure 4b, which relaxes the experimental conditions to precisely position the SPE in the cavity. For a realistic case, with a SPE that exhibits a ZPL at 670 nm that is on resonance with the cavity mode, even moderate  $Q$  values of 20,000 ( $\approx 20$  GHz linewidth), will yield a Purcell enhancement of  $\approx 530$ . To realize the high Purcell enhancement experimentally, techniques for precise positioning of SPEs to the maximum intensity of the cavity mode are required. Two different approaches can suggest solutions to this problem. The first is to find pre-existing single emitters and post-fabricate optical cavities around them. The other approach is to determin-



**Figure 4:** (a)  $Q$ -factor of the fundamental mode in a 1D photonic crystal cavity versus the imaginary refractive index coefficient of the material. (b) Spontaneous emission rate enhancement of a dipole emitter plotted as a function of distance across the 1D photonic crystal nanobeam cavity.

istically create emitters in desired locations. Recent studies of hBN single emitters show promising results for both approaches [36,37].

## Conclusion

In summary, we described and optimized a number of 2D and 1D PCC designs in free-standing and supported hBN layers. Linear defect cavities were studied as representatives of 2D photonic crystals. Period modulation of a 1D nanobeam was used to achieve a theoretical  $Q$ -factor in excess of 20,000 simply by modulating the beam width of the structure. The effect of the imaginary refractive index on the  $Q$ -factor of a nanobeam was simulated, as was the Purcell factor, showing a strong interaction between a dipole emitter and the optical mode. The designs and analyses of the hBN photonic cavities presented in this work will pave a way to a broad range of ap-

plications enabled by integrated photonic circuits based on 2D materials.

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## ORCID® IDs

Sejeong Kim - <https://orcid.org/0000-0001-9836-3608>

## References

- Caldwell, J. D.; Kretinin, A. V.; Chen, Y.; Giannini, V.; Fogler, M. M.; Francescato, Y.; Ellis, C. T.; Tischler, J. G.; Woods, C. R.; Giles, A. J.; Hong, M.; Watanabe, K.; Taniguchi, T.; Maier, S. A.; Novoselov, K. S. *Nat. Commun.* **2014**, *5*, No. 5221. doi:10.1038/ncomms6221
- Dai, S.; Ma, Q.; Andersen, T.; Mcleod, A. S.; Fei, Z.; Liu, M. K.; Wagner, M.; Watanabe, K.; Taniguchi, T.; Thiemens, M.; Keilmann, F.; Jarillo-Herrero, P.; Fogler, M. M.; Basov, D. N. *Nat. Commun.* **2015**, *6*, No. 6963. doi:10.1038/ncomms7963
- Bourrellier, R.; Meuret, S.; Tararan, A.; Stéphan, O.; Kociak, M.; Tizei, L. H.; Zobel, A. *Nano Lett.* **2016**, *16*, 4317. doi:10.1021/acs.nanolett.6b01368
- Chejanovsky, N.; Rezai, M.; Paolucci, F.; Kim, Y.; Rendler, T.; Rouabah, W.; Fávoro de Oliveira, F.; Herlinger, P.; Denisenko, A.; Yang, S.; Gerhardt, I.; Finkler, A.; Smet, J. H.; Wrachtrup, J. *Nano Lett.* **2016**, *16*, 7037. doi:10.1021/acs.nanolett.6b03268
- Feng, J.; Deschout, H.; Caneva, S.; Hofmann, S.; Lončarić, I.; Lazić, P.; Radenovic, A. *arXiv* **2017**, No. 1706.06313. <https://arxiv.org/abs/1706.06313>
- Jungwirth, N. R.; Calderon, B.; Ji, Y.; Spencer, M. G.; Flatt, M. E.; Fuchs, G. D. *Nano Lett.* **2016**, *16*, 6052. doi:10.1021/acs.nanolett.6b01987
- Kianinia, M.; Regan, B.; Tawfik, S. A.; Tran, T. T.; Ford, M. J.; Aharonovich, I.; Toth, M. *ACS Photonics* **2017**, *4*, 768. doi:10.1021/acsphotonics.7b00086
- Martínez, L. J.; Pelini, T.; Waselowski, V.; Maze, J. R.; Gil, B.; Cassabo, G.; Jacques, V. *Phys. Rev. B: Condens. Matter Mater. Phys.* **2016**, *94*, 121405. doi:10.1103/PhysRevB.94.121405
- Tran, T. T.; Bray, K.; Ford, M. J.; Toth, M.; Aharonovich, I. *Nat. Nanotechnol.* **2016**, *11*, 37. doi:10.1038/nnano.2015.242
- Tran, T. T.; Wang, D.; Xu, Z.-Q.; Yang, A.; Toth, M.; Odom, T. W.; Aharonovich, I. *Nano Lett.* **2017**, *17*, 2634. doi:10.1021/acs.nanolett.7b00444
- Englund, D.; Fushman, I.; Faraon, A.; Vučković, J. *Photonics Nanostruct. - Fundam. Appl.* **2009**, *7*, 56. doi:10.1016/j.photonics.2008.11.008
- Faraon, A.; Santori, C.; Huang, Z.; Acosta, V. M.; Beausoleil, R. G. *Phys. Rev. Lett.* **2012**, *109*, 033604. doi:10.1103/PhysRevLett.109.033604
- Kuhn, A.; Ljunggren, D. *Contemp. Phys.* **2010**, *51*, 289. doi:10.1080/00107511003602990
- Müller, K.; Fischer, K. A.; Rundquist, A.; Dory, C.; Lagoudakis, K. G.; Sarmiento, T.; Kelaita, Y. A.; Borish, V.; Vučković, J. *Phys. Rev. X* **2015**, *5*, 031006. doi:10.1103/PhysRevX.5.031006
- Sato, Y.; Tanaka, Y.; Upham, J.; Takahashi, Y.; Asano, T.; Noda, S. *Nat. Photonics* **2012**, *6*, 56. doi:10.1038/nphoton.2011.286
- Benson, O. *Nature* **2011**, *480*, 193. doi:10.1038/nature10610
- Harris, N. C.; Bunandar, D.; Pant, M.; Steinbrecher, G. R.; Mower, J.; Prabhu, M.; Baehr-Jones, T.; Hochberg, M.; Englund, D. *Nanophotonics* **2016**, *5*, 456. doi:10.1515/nanoph-2015-0146
- Santori, C.; Fattal, D.; Vučković, J.; Solomon, G. S.; Yamamoto, Y. *Nature* **2002**, *419*, 594. doi:10.1038/nature01086
- Bracher, D. O.; Zhang, X.; Hu, E. L. *Proc. Natl. Acad. Sci. U. S. A.* **2017**, *114*, 4060. doi:10.1073/pnas.1704219114
- Faraon, A.; Barclay, P. E.; Santori, C.; Fu, K.-M. C.; Beausoleil, R. G. *Nat. Photonics* **2011**, *5*, 301. doi:10.1038/nphoton.2011.52
- Cassabo, G.; Valvin, P.; Gil, B. *Nat. Photonics* **2016**, *10*, 262. doi:10.1038/nphoton.2015.277
- Lin, Y.; Connell, J. W. *Nanoscale* **2012**, *4*, 6908. doi:10.1039/c2nr32201c
- Liu, Z.; Gong, Y.; Zhou, W.; Ma, L.; Yu, J.; Idrobo, J. C.; Jung, J.; MacDonald, A. H.; Vajtai, R.; Lou, J.; Ajayan, P. M. *Nat. Commun.* **2013**, *4*, 2541. doi:10.1038/ncomms3541
- Tran, T. T.; Elbadawi, C.; Totonjian, D.; Lobo, C. J.; Grosso, G.; Moon, H.; Englund, D. R.; Ford, M. J.; Aharonovich, I.; Toth, M. *ACS Nano* **2016**, *10*, 7331. doi:10.1021/acsnano.6b03602
- Elbadawi, C.; Tran, T. T.; Kolibal, M.; Šikola, T.; Scott, J.; Cai, Q.; Li, L. H.; Taniguchi, T.; Watanabe, K.; Toth, M.; Aharonovich, I.; Lobo, C. *Nanoscale* **2016**, *8*, 16182. doi:10.1039/C6NR04959A
- El Kurdi, M.; Checoury, X.; David, S.; Ngo, T.; Zerounian, N.; Boucaud, P.; Kermarrec, O.; Campidelli, Y.; Bensahel, D. *Opt. Express* **2008**, *16*, 8780. doi:10.1364/OE.16.008780
- Schwagmann, A.; Kalliakos, S.; Ellis, D. J. P.; Farrer, I.; Griffiths, J. P.; Jones, G. A. C.; Ritchie, D. A.; Shields, A. J. *Opt. Express* **2012**, *20*, 28614. doi:10.1364/OE.20.028614
- Akahane, Y.; Asano, T.; Song, B.-S.; Noda, S. *Nature* **2003**, *425*, 944. doi:10.1038/nature02063
- Kim, S.-H.; Kim, G.-H.; Kim, S.-K.; Park, H.-G.; Lee, Y.-H.; Kim, S.-B. *J. Appl. Phys.* **2004**, *95*, 411. doi:10.1063/1.1633645
- Hunger, D.; Steinmetz, T.; Colombe, Y.; Deutsch, C.; Hänsch, T. W.; Reichel, J. *New J. Phys.* **2010**, *12*, 065038. doi:10.1088/1367-2630/12/6/065038
- Quan, Q.; Burgess, I. B.; Tang, S. K. Y.; Floyd, D. L.; Loncar, M. *Opt. Express* **2011**, *19*, 22191. doi:10.1364/OE.19.022191
- Kim, S.; Kim, H.-M.; Lee, Y.-H. *Opt. Lett.* **2015**, *40*, 5351. doi:10.1364/OL.40.005351
- Kim, S.; Ko, H.; Lee, C.; Kim, M.; Kim, K. S.; Lee, Y.-H.; Shin, K.; Cho, Y.-H. *Adv. Mater.* **2016**, *28*, 9765. doi:10.1002/adma.201603368
- Kim, S.; Ahn, B.-H.; Kim, J.-Y.; Jeong, K.-Y.; Kim, K. S.; Lee, Y.-H. *Opt. Express* **2011**, *19*, 24055. doi:10.1364/OE.19.024055
- Kreuzer, C.; Riedrich-Möller, J.; Neu, E.; Becher, C. *Opt. Express* **2008**, *16*, 1632. doi:10.1364/OE.16.001632
- Kianinia, M.; Bradac, C.; Wang, F.; Sontheimer, B.; Tran, T. T.; Nguyen, M.; Kim, S.; Xu, Z.-Q.; Jin, D.; Schell, A. W.; Lobo, C. J.; Aharonovich, I.; Toth, M. *arXiv* **2017**, No. 1709.08683. <https://arxiv.org/abs/1709.08683>

37. Xu, Z.-Q.; Elbadawi, C.; Tran, T. T.; Kianinia, M.; Li, X.; Liu, D.; Hoffman, T. B.; Nguyen, M.; Kim, S.; Edgar, J. H.; Wu, X.; Song, L.; Ali, S.; Ford, M. J.; Toth, M.; Aharonovich, I. *arXiv* 2017, No. 1710.07010.  
<https://arxiv.org/abs/1710.07010>

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