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# Compact Multilayer Bandpass Filter Using Low-Temperature Additively Manufacturing Solution

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Abstract-This presented additively paper an manufactured bandpass filter (BPF) based on а second-order stub-loaded resonator consisting of multi-metal layer components. The proposed BPF is fabricated by a low-temperature (140°) additively manufactured electronics (AME) solution that can fabricate conductive and dielectric materials simultaneously with multi-metal-layer and flexible interlayer distance. By reducing the interlayer distance, constant inductance and capacitance can be realized in smaller sizes, which helps to achieve device minimization. Taking advantage of this inkjet printing technology, a second-order multi-metal layer resonator is proposed. To understand the principle of the BPF, an equivalent circuit with odd- and even-mode analysis is demonstrated. For verification, the frequency response of the circuit's mathematical model is calculated to compare with the electromagnetic simulation results. Good agreement can be achieved among the calculated, simulated, and measured results. The proposed BPF is designed at 12.25 GHz with a bandwidth of 40.8% and a compact size of 2.7 mm × 1.425 mm × 0.585 mm or 0.186λg × 0.098 $\lambda$ g× 0.040 $\lambda$ g, which is suitable for circuit-in-package applications in television programs, radar detection and satellite communications.

*Index Terms*—bandpass filter (BPF), wide stopband, multi-metal-layer components, additively manufactured electronics (AME), low-temperature, 3D inkjet printing.

#### I. Introduction

ITH the rapid development of additive manufacturing technology, more and more multi-material 3D printers can fabricate conductive and dielectric materials simultaneously, which have demonstrated their capabilities in electronic devices integration and 3D packaging [1]-[5]. However, it is still a challenge for these 3D printers to fabricate

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Fig. 1. The proposed low-temperature additively manufacturing electronics (AME) process.

seamless multi-metal-layer in a single substrate. According to the open literature, multi-metal-layer technologies have an outstanding advantage in device minimization such as low-temperature co-fired ceramic (LTCC) [6]-[10] or semiconductor technologies [11]-[13]. For example, reference [8] proposes a compact LTCC diplexer with high-isolation. In [11], an on-chip miniaturized BPF is presented, which is based on a grounded center-tapped ring resonator with shunt capacitive loading. Although both designs demonstrated device minimization, these technologies are expensive and require complicated fabrication processes and longer production time. As an emerging technology, low-temperature additively manufacturing technologies have been applied in integrated electronic devices and packaging, such as electronic components [14], antenna arrays [15], terahertz lens [16], redistribution layers [17], combiner [18], transistors [19], and enzymatic biofuel [20].

In this paper, a low-temperature additively manufactured electronics (AME) solution is introduced for bandpass filter (BPF) design, taking advantage of dual-materials (conductive and dielectric materials) printing, as shown in Fig. 1. Nano Dimension's DragonFly LDM system is used for AME designs with 4pL printer heads. This system simultaneously jets both dielectric and conductive inks, hence simultaneously creating multilayer conductive structures within a single dielectric substrate. Compared with LTCC and semiconductor technologies, the AME solution is flexible on electronic device fabrication with design freedom adjusting the interlayer



Fig. 2 Simulated inductances and quality factors of multilayer spiral inductors with different dielectric-layer thickness h.



Fig. 3 Comparison of simulated capacitances of multilayer capacitors with different dielectric-layer thickness t.

distance. The printing process is carried out at a typical temperature of 140°C in a standard ambient atmosphere. For proof-of-concept, a multilayer BPF with vertically integrated capacitors and inductors has been designed, analyzed, fabricated and measured. Finite element electromagnetic field simulation analysis software, ANSYS High-Frequency Structure Simulator (HFSS), is used in the simulation, and the Vector Network Analyzer (VNA) used in the measurement is a Keysight PNA N5225B.

# II. ADDITIVELY MANUFACTURED COMPONENTS FOR BANDPASS FILTER MINIATURISATION

The printed BPF is composed of the inductance and capacitance components. The proposed AME technique has the advantage of circuit miniaturization, which is demonstrated by using an inductor and a capacitor, as the examples.

#### A. Miniaturized 3D spiral inductor using AME process.

As shown in Fig. 1, the fabrication procedure of a 3D spiral inductor is demonstrated using the proposed low-temperature AME process. Silver nanoparticles and acrylate inks can be printed simultaneously in this AME process. During the manufacturing process, for *Step 1*, an ultraviolet (UV) lamp with a wavelength of 395 nm is used for curing the acrylate inks after they are ejected out from the nozzles. In *Step 2*, a metal

strip line is printed on the acrylate layer. A near-infrared radiation (NIR) lamp with a wavelength of 0.75-1.4  $\mu$ m and 140°C to 170°C process temperature is used to sinter the conductive inks. In *Steps 3 – 6*, acrylate and conductive inks are printed layer by layer to construct the 3D spiral inductor. Vias can be printed simultaneously in the acrylate layer to connect metal strip lines in different layers, as depicted in *Step 3*. Where a via is desired, no dielectric ink is deposited. This via can remain a through via hole or can be filled with the conductive silver nanoparticle ink. The fabrication is done layer by layer. Hence the silver is in the area where the filled via is being built, and the dielectric is around it. Taking advantage of the AME solution, designing 3D components and devices would be flexible.

For vertical spiral inductors, the inductance can be calculated by

$$L = K \frac{\mu_0 \mu_r N^2 A_e}{H} \tag{1}$$

where  $\mathcal{H}_0$  is the permeability of free space,  $\mathcal{H}_i$  is the relative permeability of the material, N is the total number of turns; H is the total height of coil,  $A_e$  is the section area of the coil, and K is a coefficient that depends on the ratio of r (radius of the coil) to H. To investigate the relationship between the L and H, Fig. 2 demonstrates the comparison of simulated inductances with different dielectric-layer thickness h. When h drops from 0.15 mm to 0.05 mm, the inductance increases gradually. In addition, Fig. 2 shows the quality factor versus different values of h. With the decline of h, the inductance density increases, which means that a higher inductance value can be obtained with a lower profile and constant plane size. A drop in peak quality factor is observed because increasing inductance usually involves increasing magnetic layer thickness, which leads to higher eddy currents.

# B. Miniaturized capacitors with reduced interlayer distance

According to reference [21], for a multilayer capacitor, the capacitance "C" is given in the equation below:

$$c = \frac{\varepsilon_0 \varepsilon_s S(n-1)}{t} \tag{2}$$

where  $\varepsilon_0$  is the dielectric constant of free space,  $\varepsilon_s$  is the relative dielectric constant of the material, *S* is the effective area of the inner electrode. *n* is the number of inner electrodes, and *t* is the thickness of the dielectric layer.

Obviously, with the increase of the number of inner electrodes or raising the inner electrode effective area, the capacitance will grow. That is also the most common method to design a capacitor in previous literature. However, in this paper, the effect of the thickness of the dielectric layer are discussed. As shown in Fig. 3, with the decrease of dielectric-layer thickness *t* from 0.15mm to 0.10mm, capacitance increases significantly while the capacitor size keeps unchanged, consistent with equation (2). In other words, the adjustable interlayer distance can contribute to a smaller plane size and lower profile.

# III. MULTILAYER BPF AND ITS EQUIVALENT CIRCUIT

# A. Bandpass filter configuration

In Fig. 4(a), the 3-D view of the proposed BPF is presented.



Fig. 4 (a) 3-D view and (b) Side view, (c) Layout of each layer, (d) Photograph of the 3D printed BPF with testing accessories and structural details of the filter under microscope.

Taking advantage of the proposed AME solution, the number of metal layers and the dielectric layer thickness among metal layers can be flexibly defined for circuit design. In this work, as shown in Fig. 4(b), five metal layers are employed to design the BPF, where layer 5 is used as the ground. The constructions of feeding lines and resonators are distributed from layer 1 to layer 4. The resonant part of the BPF consists of two  $\lambda/4$  resonators, where three metal layers are stacked to form a compact capacitor. Two stepped-impedance stubs, located on layer 1 and layer 3, respectively, are loaded at the one third point to improve the of out-of-band suppression. Two up-raised grounds located at layer 3 are connected to the RF ground by 24 vias, respectively, to adjust the impedance match. Fig. 4(c) depicts the layout of the metal layers. Parameters of the proposed BPF are determined as follows (all in mm):  $D_1=3.3$ ,  $D_2=1.6$ ,  $D_3=1.28$ ,  $D_4=1.3$ ,  $D_5=1.305$ ,  $D_6=0.85$ ,  $D_7=0.4$ ,  $D_8=0.5$ ,  $D_9=0.85$ ,  $D_{10}=0.425$ ,  $D_{11}=0.7$ ,  $W_1=0.5$ ,  $W_2=0.125$ ,  $W_3=0.4$ ,  $W_4=0.65$ ,  $W_5=0.15$ ,  $W_6=0.4$ ,  $r_1=0.1$ ,  $g_1=0.7$ . The dimension of BPF excluding feeding lines is 2.7 mm × 1.425 mm × 0.585 mm or 0.186  $\lambda g \times 0.098 \lambda g \times 0.040 \lambda g$ , where  $\lambda g$  is the guided wavelength at the center frequency. Fig. 4(d) shows the photograph of an additively manufactured prototype with testing accessories and structural details of the filter under a microscope.

#### B. Equivalent lumped component circuit

To facilitate the analysis of the proposed multi-metal layers BPF, an equivalent circuit based on the unfolded layout of the BPF is presented and discussed. Fig. 5 presents the equivalent lumped component (LC)-circuit model of the presented EM structure, where inductors and capacitors represent the high-impedance lines and the capacitive couplings, respectively. L<sub>Fi</sub> (i=1, 2) denotes the self-inductances of the feeding lines.  $C_{\rm F}$  is the capacitive coupling between the soldering area of SMK connector and the ground. C<sub>b</sub> represents the coupling capacitance between the feeding line and the resonator.  $C_p$  is the bypass capacitance of the feeding line to the ground. In addition, the source to load coupling is modeled by  $C_{\text{S-L}}$ . The inductors and capacitors of the resonators are  $L_1$  and  $C_1$ , where  $L_1$  and  $C_1$  represent the parasitic inductor and the capacitance of this three-dimension interdigital capacitor, respectively. L2 is the self-inductance of the quarter wavelength resonator. The loaded stepped-impedance is represented by  $C_2$ ,  $L_3$  and  $C_3$ ,  $L_4$ .  $C_2$ ,  $L_3$  denotes the low-impedance segment. The  $C_3$  and  $L_4$  denote the high-impedance segment.

# C. Analysis of equivalent LC-circuit model

For circuit analysis, an equivalent *LC*-circuit of resonators is depicted in Fig. 6(a). Since the two resonators are symmetrical about AA', odd- and even-mode circuits can be obtained, as shown in Figs. 6(b) and (c). The loop equation method is used for circuit analysis, of which the loop currents  $I_i$  (*i*=1, 2, 3, 4) are marked in Figs. 6(b) and (c). The expressions of input admittance are listed to calculate the resonant frequencies of transmission poles (TPs) and transmission zeros (TZs).

Applying Kirchhoff's voltage law (KVL) to the odd-mode equivalent circuit in Fig. 6(b), the voltage equations of the four loops are expressed as:

$$\frac{I_1}{sC_1} + sL_1I_1 + sL_2(I_1 - I_2) - U_{\text{odd}} = 0$$
(3)

$$\frac{(I_2 - I_3)}{sC_2} + sL_2(I_2 - I_1) = 0 \tag{4}$$

$$\frac{(I_3 - I_2)}{sC_2} + sL_3I_3 + \frac{(I_3 - I_4)}{sC_3} = 0$$
(5)

$$sL_4I_4 + \frac{(I_4 - I_3)}{sC_3} = 0 \tag{6}$$

Similarity, applying KVL to the even-mode equivalent circuit in Fig. 6(c), the voltage equations of the three loops are



Fig. 5 Equivalent LC-circuit model based on unfolded layout of the proposed BPF (Derived from Fig. 4(a) for illustration of the LC components).

expressed as:

$$\frac{I_1}{sC_1} + sL_1I_1 + sL_2(I_1 - I_2) - U_{\text{even}} = 0$$
<sup>(7)</sup>

$$\frac{(I_2 - I_3)}{sC_2} + sL_2(I_2 - I_1) = 0$$
(8)

$$\frac{(I_3 - I_2)}{sC_2} + sL_3I_3 + \frac{I_3}{sC_3} = 0$$
<sup>(9)</sup>

where  $U_{odd}$  and  $U_{even}$  denote the input voltages of odd- and even- mode circuits, respectively. *s* is a complex variable in proportion to angular frequency  $\omega$ . Therefore, the relationships between input voltages and input admittances in Figs. 6(b) and (c) can be written as:

$$Y_{odd}(s) = \frac{I_{1\_odd}(s)}{U_{odd}(s)} = \frac{I_{1\_odd}(j\omega_{odd})}{U_{odd}(j\omega_{odd})}$$
(10)

$$Y_{even}(s) = \frac{I_{I\_even}(s)}{U_{even}(s)} = \frac{I_{I\_even}(j\omega_{even})}{U_{even}(j\omega_{even})}$$
(11)

where  $\omega$  denotes angular frequency, and *j* is a pure imaginary unit number. Using MATLAB, the expression of input admittances consists of complex variable s and lumped LCelements can be deduced by solving equations (2)-(11). Since  $\omega = 2\pi f$ , the frequencies  $f_{odd}$  and  $f_{even}$  (corresponding to odd- and even-mode) can be calculated when  $Y_{odd}(j\omega) = 0$  and  $Y_{even}(j\omega) =$ 0, respectively. Meanwhile, the TZs can be determined when  $Y_{\text{odd}}(j\omega) = Y_{\text{even}}(j\omega)$ . For verification, the equivalent LC-circuit is optimized to match the finite element electromagnetic field simulation result of the EM structure. As shown in Fig. 7(a), the simulated results of equivalent LC-circuit carried by ADS are illustrated. Good agreement can be obtained between the calculated results of LC-circuit and simulated results of the EM structure. Substituting the obtained values of lumped elements in equivalent circuit:  $L_1 = 0.016 \text{ nH}, L_2 = 0.589 \text{ nH}, L_3 = 0.725 \text{ nH},$ L<sub>4</sub>=1.247 nH, C<sub>1</sub>=0.345 pF, C<sub>2</sub>=0.18 pF, C<sub>3</sub>=0.11 pF into equations (12)-(13), one TZs  $f_{zi}$  (i=1) and two TPs  $f_{pi}$  (i=1,2) can be obtained in the frequency range of 0-35 GHz through the calculations. Finally, the estimated TZs and TPs are marked at Fig. 7(a) as well. Compared with the simulated results of TZs and TPs:  $f_{z1}$  =19.85 GHz,  $f_{p1}$  =11.31 GHz and  $f_{p2}$  =13.18 GHz, the calculated results agree well with the simulated ones. In addition, to improve the skirt selectivity, (source-load) S-L coupling is induced by adding a capacitor between two ports, as shown in Fig. 7(b). The locations of the TZs are dependent on the mixed electromagnetic (EM) coupling coefficient between the two resonators and the (source-load) S-L coupling construction. Since these two resonators of the BPF are



Fig. 6 (a) Equivalent circuit of the resonators and its (b) Odd- and (c) Even-mode equivalent circuits.

connected by via, which results that the magnetic coupling is larger than the electric one. Therefore, a TZ would be excited at the upper (magnetic-dominant coupling) stopband (TZ<sub>2</sub>). As shown in Fig. 8(a), the frequency response of the coupled resonators against spacing  $D_7$  is depicted, which shows the variation of the EM coupling between the two resonators. Lower-band TZ (TZ<sub>1</sub>) and upper-band TZ (TZ<sub>2</sub>) are marked in the picture. By tuning  $D_7$ , the location of the TZ<sub>2</sub> can be adjusted at the upper stopband, which significantly enhances the out-of-band attenuation of the proposed filter, while the TZ<sub>1</sub> keeps unchanged.

In addition, the simulated frequency response against  $g_1$  is given for the TZ created in the low-band (TZ<sub>1</sub>), as shown in Fig. 8(b). When  $g_1$  decreases from 0.3 to 0.1 mm, the TZ<sub>1</sub> increases from 5.88 to 8.37 GHz, whereas the TZ<sub>2</sub> changes slightly without affecting the filter bandwidth. The closer the TZs to the passband, the better the frequency selectivity can be achieved.

For the frequency response of the TPs, by solving  $Y_{odd}(j\omega) = 0$  and  $Y_{even}(j\omega) = 0$ , it can be found that the even-mode resonant frequency  $f_{even}$  is the function of  $L_4$ , while  $f_{odd}$  is the function of both  $L_2$  and  $L_4$ .  $L_2$  and  $L_4$  correspond to the length of  $D_4$  and  $D_{10}$  in the EM structure, respectively. For validation, the simulated responses against  $D_4$  and  $D_{10}$  are shown in Figs. 8(c) and (d). When  $D_4$  rises from 0.9 mm to 1.3 mm,  $f_{p1}$  keeps constant while  $f_{p2}$  decreases gradually. Besides, with the development of  $D_{10}$ ,  $f_{p1}$  and  $f_{p2}$  decline simultaneously. The variations are consistent with the analysis and verify that the bandwidth and center frequency can be flexibly designed.



Fig. 7 Simulated frequency responses of (a) the BPF (by HFSS) as well as the equivalent *LC*-circuit in Fig. 6 (by ADS) and (b) Simulated frequency responses of the BPF with or without S-L coupling.

The unloaded quality factor of the resonators is important for filter performance. In Fig. 9(a), for a single resonator used in the proposed BPF (center frequency = 12.25 GHz), the simulated unloaded Q-factor versus the length  $L=D_9+D_{10}+D_{11}$ is presented. According to Fig.9(a), when L increases, the unloaded  $Q_u$  decreases due to dielectric and conductor losses. Besides, the external quality factor  $Q_e$  with varied coupling gap G is depicted in Fig. R2 (b). When G increases from 0.05 to 0.5mm, a higher  $Q_e$  of the filter can be obtained, which causes a decrease of the bandwidth.

#### D. Suppression of third harmonic and wide stopband

Fig. 10(a) shows the structure of the quarter-wavelength resonator used in the proposed BPF. A short-circuited stepped impedance stub is loaded at the trisection-point of this resonator closing to the input terminal. The electrical lengths of the quarter-wavelength transmission line are referred to  $\theta_1$  and  $2\theta_1$ , respectively. The short-circuited stub position can be defined by analyzing the voltage distributions of the fundamental mode and the third harmonic to improve the out-of-band suppression. It is obvious that the shorted terminal of the quarter-wavelength resonator is the common voltage zero point of the fundamental mode and the third harmonic mode, while the trisection-points near the input terminal is the zero-voltage point only for third harmonic mode. When the shorted stub is located at the zero-voltage point of the third harmonic mode, the third



Fig. 8 Simulated frequency responses against (a)  $D_7$ , (b)  $g_1$  with coupling topology inserted, (c)  $D_4$  and (d)  $D_{10}$ .



Fig. 9 (a) Simulated Unloaded  $Q_u$  of the proposed single resonator against a different *L* (b) and extracted external quality factors  $Q_e$  against a different coupling gap *G*.

harmonic mode signal will not get through the resonator, as shown in Fig. 10(b). Thus, out-of-band suppression can be developed by adjusting the location of the loaded stub. In Fig. 10(c), with the decrease of  $D_9$  from 1.3 mm to 1 mm, the loaded stub approaching the zero-voltage point of the third harmonic mode gradually, and the out-of-band suppression is improved from 12.5 dB to 27.3 dB.

#### IV. MEASUREMENT RESULT

The performance of the designed BPF is measured from 1 GHz up to 35 GHz. Good agreement between the simulated and measured results is observed in Fig. 11. The dielectric constant of the additive manufactured substrate is 2.75, and the loss tangent is 0.025. [22-23]. The conductivity of the silver film is  $3.15 \times 10^6 - 2.52 \times 10^7 \sigma$  (S/m) at 20 °C. As demonstrated in Fig. 11, compared with the simulated results achieved by HFSS, the measured insertion loss and return loss are 3.42 dB and 16.65 dB at the center of 12.25 GHz, respectively. In the measurement, the connection between the SMK connector and the filter will cause inter-connection loss, which will lead to a

Ref.	f <sub>0</sub> (GHz)	*FBW (%)	Number of TZs	Filter order	Stopband (GHz)	$\varepsilon_{ m r}/{ m tan}\delta$	Size (mm <sup>2</sup> )/ $(\lambda_g^2)$	Thickness $(mm)/(\lambda_g)$	Fabricated process
[2]	1.6	22	0	4	$0.25 f_0 (20 \text{dB})$	2.75/0.015	20×83/0.177×0.735	3.6/0.032	Additive
	2.45	12	0	4	$0.16 f_0 (20 \text{dB})$	2.35/0.02	11.15×50/0.140×0.626	3.6/0.045	manufacturing
[11]	33	42.4	0	1	$1.58 f_0 (20 dB)$	N.A.	0.11×0.28/N.A.	N.A.	0.13-µm (Bi)-CMOS
[12]	59.5	21.68	2	2	N.A.	N.A.	0.240×0.225/0.144×0.135	N.A.	0.18-µm CMOS
[24]	4.3	44.2	1	3	$2.6 f_0 (20 \text{dB})$	5.9/N.A.	12×6/0.418×0.0209	1/0.035	LTCC
This work	12.25	40.8	2	2	1.9 f <sub>0</sub> (20dB)	2.75/0.025	2.7×1.425/0.186×0.098	0.508/0.04	AME solution

TABLE I Comparison with Some Other BPF

\*FBW= fractional Bandwidth



Fig. 10 (a) Voltage distribution of the quarter-wavelength resonator used in the proposed BPF (b) Electric filed distribution of the proposed filter at the resonate frequency of third harmonic mode when (1)  $D_9=1.4$ mm and (2)  $D_9=1$  mm and (c) Simulated  $|S_{21}|$  against  $D_9$ .

deteriorated insertion loss. At the same time, the error of the thickness in the manufacturing process is also the reason for a deteriorated insertion loss. A 3-dB bandwidth is achieved from 9.58 GHz to 14.57 GHz with the FBW of 40.8%. Two TZs located at 3.51 and 18.51 GHz that can improve the selectivity are generated. In addition, due to the loaded short-circuited stepped impedance stub the upper stopband with more than 20



Fig. 11 Simulated and measured frequency response with a zoom-in view of the passband insert loss of the AME printed BPF.

dB rejection level is extended to 23.24 GHz (from 16.65 to 39.89 GHz) or 1.90  $f_0$ . As can be seen, the designed BPF working at 12.25 GHz has the merits of multiple transmission zeros and poles, a compact size, and a wide stopband that can be assigned to television programs or be utilized in radar detection and satellite communications.

Table I shows the comparison of the proposed BPF with other state-of-the-art BPFs fabricated with different fabricated processes. Compared with [2], the BPF proposed in our work has a larger FBW, two flexibly designed transmission zeros, a wider stopband, and a smaller size. In addition, compared with traditional multilayer processes, such as CMOS and LTCC technology, the proposed BPF has a good FBW, high-band and low-band TZs, a good stopband, and a competitive footprint. Furthermore, AME solution used in this paper has the advantages of a low-cost and short processing cycle.

### V. CONCLUSION

In this paper, a low-temperature additively manufactured electronics (AME) technology for multilayer 3D BPF design is presented. Equivalent circuits and Kirchhoff's voltage law are used for analyzing the frequency response of the printed BPF. The proposed additive manufacturing process functions at a typical temperature of 140°C in a standard ambient atmosphere. Good agreement among the calculated, EM simulated and measured results validate the feasibility of the proposed AME solution for miniaturized bandpass filter design. The proposed AME solution has great potential in circuit-in-package applications.

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