Natural Ventilation Induced by Solar Chimneys

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Abstract
Natural ventilation flow through a two-dimensional but real-sized square room is investigated numerically, using a commercial Computational Fluid Dynamics (CFD) software package. The flow is induced by a solar chimney positioned on the room's roof, and it is desired to have this flow passing through the lower part of the room for ventilation purpose. The chimney in turn is in the form of a parallel channel with one plate kept at a uniform temperature that is higher than that of the ambient air (by up to 40°C), while the other plate and all of the room's walls are insulated. Effects on ventilation flow rate and flow pattern due to a range of changing factors are investigated.

Introduction
People spend most of their time indoors. A comfortable indoor environment is thus essential for the occupants' good health and productivity. Buildings are responsible for about half of a modern society's total energy consumption. HVAC (Heating, Ventilation and Air-Conditioning), in turn, accounts for a major proportion of this energy demand, thus estimated to be about 68% in non-industrial (commercial and residential) buildings [16]; HVAC is often used to provide thermal comfort to the occupants. Minimising HVAC energy consumption will thus result in great economic benefits. It also contributes beneficially to the issue of sustainable future and climate change, by reducing fuel burning.

Natural ventilation can be used to help reduce significantly HVAC energy demand. A solar chimney (thermal chimney) helps enhance natural ventilation by being heated with solar radiation, causing hot air to rise and inducing the ventilation flow. This device has been much investigated [1-6, 9-10, 12-17]; but there are still many factors affecting its performance (measured by the induced flow rate, for example) not yet considered, especially in building settings. This work investigates computationally natural ventilation induced by roof-mounted solar chimneys through a two-dimensional but real-sized square room. Air-mass-flow rate and flow pattern will be considered in terms of chimney-room configuration, the presence of a vertical partition in the room, chimney's length and inclination, position of the chimney's heated plate, and its temperature.

Modelling and Computation
The flow model is depicted in figure 1. A two-dimensional but real-sized square room of width W = 3m and height H = 3m is considered. Flow is induced by a solar chimney (SC) positioned on the room's roof. The chimney in turn is in the form of a parallel channel with one plate kept at a uniform temperature that is higher than the ambient-air temperature which is assumed to be fixed at 300K (27°C). Two inclination states of the chimney are considered: "inclined" when the chimney is inclined at an angle of 30° above the horizontal (shown in figure 1), and "vertical" when chimney is vertical. When the chimney is inclined, its heated plate can be either on the left or right, while if the chimney is vertical, its heated plate can be on the left or right. Chimney's length is 1.56m in most cases; but a length twice as long, at 3.12m, is also considered.

Air is induced into the chimney through its lower opening which is a slot on the room's roof either at B (the middle) or C (right corner), and flows out of the chimney through its upper opening. Air enters the room via one or 2 slots located at A (on the roof), D (at bottom of the left vertical wall) or E (bottom of right wall). Opening slots' locations A, B, C, D, and E are shown in figure 1; all slots are 0.3m wide. Note that when SC is inclined, the distance between its 2 plates is only 0.260m (-30°C) but when SC is vertical, this distance is 0.3m. Another feature is the presence or absence of a vertical partition which is a thin solid wall hanging down from the roof, ending at 0.3m above the floor. Table 1 shows the various configurations (series) considered. Computational flow domain thus consists of the room's space plus the space between the SC's plates.

All fluid properties are assumed to be constant and corresponding to those of air at 300K (constant ambient temperature T_A) and standard pressure at sea level (101.3kPa); but Boussinesq approximation is also assumed for the buoyancy force arising from density variation as a result of temperature change. The following values of molecular properties are used (using common notation): ρ = 1.161 kg/m^3; μ = 1.846x10^-5 Ns/m^2; ν = 1.589x10^-5 m^2/s; k = 0.0263 W/m-K; α = 1007.17 kg/J/K; α = 2.25x10^-3 m^2/s; Pr = ρa/μ = 0.707; β = 1/T_A = 1/300 K^-1. With these values, the Rayleigh number Ra, which is a key parameter in free convection and based on the inclined chimney's internal gap size (0.260m), is Ra = 1.60x10^4 ΔT, where ΔT is the temperature difference between the SC's heated plate and the ambient air. In this work, ΔT can be up to 40°C. Since turbulence is expected at Ra > 10^6 for free convection in similar configurations [11], all cases considered in this work are taken to be turbulent. Following a common practice, turbulent Prandtl number is taken to be constant at 0.9. It should be noted, however, that since laminar flow is a special case of turbulent flow with zero turbulence level, use of a turbulence model (see below) in laminar flow should not affect the computational results' correctness. Thus, for example, with a case from the 1 series with heated plate's temperature 301K wherein laminar flow is expected in the room, computation with a turbulence model (versus laminar-flow computation) results in a difference of only 0.02% in the air-mass-flow rate through each of the room's inlets.

A Reynolds-Averaged Navier-Stokes (RANS) formulation is used, wherein turbulence affects the mean flow through a turbulent viscosity μ_t. Turbulent stresses are assumed to be
The turbulence-model variables $K$ and $s$ are prescribed as follows:

Referring to figure 1, boundary conditions for the mean variables (velocity components, pressure and temperature) are as follows:

- All openings on the flow domain's boundary (SC's upper opening, inlets to the room at A (figure 1), and D and E in some configurations) are prescribed with ambient conditions; the fluid has constant ambient pressure and temperature, namely $p = 0$ (gauge, without the hydrostatic component), $T = T_a = 300K$. However, the thermal condition here applies only on those sections of an opening where there is inflow; if the computation reveals inflow on any sections, the constant temperature condition there will be ignored; instead, temperature will be computed instead.

- Isothermal wall condition is prescribed to the heated plate of the chimney: zero velocity, $T - T_{plate}$ (constant, and $> T_a$)

- All other solid surfaces (walls, roof and floor of the room, SC's non-heated plate, and the partition when it is present) are adiabatic walls: zero velocity, $\frac{\partial T}{\partial n} = 0$

The turbulence-model variables $K$ and $s$ are prescribed as follows:

- On all solid surfaces: default solid-surface condition of the software package (see below) is adopted; this entails $K = 0$ and $s = 0$, following Chien [7]

- At all openings, $K$ and $s$ are assumed to be constant. However, these conditions apply only on those sections of an opening where there is inflow; if the computation reveals inflow on any sections, the prescribed values for $K$ and $s$ there will be ignored; instead, these will be computed. It is found that the small turbulence level prescribed on the openings has very small effects on the results. For example, when an arbitrary value of $K = 1 \times 10^{-5} m^2/s^3$ and $s = 1 \times 10^{-5} m^2/s^3$ are prescribed, the air-mass-flow rate varies by less than 0.3% from when $K$ and $s$ are both zero. From considerations similar to this, $K = 1 \times 10^{-5} m^2/s^3$ and $s = 1 \times 10^{-5} m^2/s^3$ are used on all openings.

<table>
<thead>
<tr>
<th>Series</th>
<th>Partition</th>
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<th>SC</th>
<th>SC Inclination</th>
<th>SC's Heated Plate</th>
</tr>
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<tbody>
<tr>
<td>A</td>
<td>Yes</td>
<td>A</td>
<td>C</td>
<td>inclined</td>
<td>Lower</td>
</tr>
<tr>
<td>B</td>
<td>No</td>
<td>A</td>
<td>C</td>
<td>inclined</td>
<td>Lower</td>
</tr>
<tr>
<td>C</td>
<td>No</td>
<td>D</td>
<td>C</td>
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</tr>
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<tr>
<td>E</td>
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<td>D &amp; E</td>
<td>C</td>
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<td>Lower</td>
</tr>
<tr>
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</tr>
<tr>
<td>H</td>
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<td>D &amp; E</td>
<td>B</td>
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</tr>
<tr>
<td>I</td>
<td>No</td>
<td>D</td>
<td>B</td>
<td>inclined</td>
<td>Lower</td>
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<tr>
<td>J</td>
<td>No</td>
<td>E</td>
<td>C</td>
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<td>Upper</td>
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<tr>
<td>K</td>
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<td>C</td>
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<td>Lower</td>
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<tr>
<td>N</td>
<td>No</td>
<td>A</td>
<td>C</td>
<td>inclined</td>
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</table>

Table 1. Series (configuration) of the cases considered. In each series temperature of the solar chimney (SC)’s heated plate is varied. SC’s length is 1.56m, except for series N wherein this length is 3.12m.

The commercial Computational Fluid Dynamics (CFD) software package CFD-ACE from the ESI Group is used for the computation. The package is quite well known, and its validation is thus assumed to have been adequate. Some further tests had also added positively to its validity [8]. Numerical scheme is the Finite Volume method, and the coupled system of governing equations is solved iteratively for the two mean velocity components, mean temperature and pressure, plus $K$ and $s$. Both upwind and central-difference differencing schemes are used. A convergence criterion of reduction of residuals in the solved variables by 3 orders of magnitude is adopted. This is adequate; comparison of the solutions with residual reduction of 3 orders of magnitude and those with 4 orders of magnitude shows very small difference. For example, with a test case from series A with plate temperature 310K, the above change of convergence criterion results in the air-mass-flow rate changing by only 0.0013%. Computation is done with 64-bit precision.

Grid convergence tests have also been performed to ascertain the adequacy of the grid patterns used. For example, with a case from t series with plate temperature 320K, as the number of grid points on each edge (see figure 1) is increased by 20% (resulting in 44% increase in the number of 2-D computational cells), change in the net mass in-flow to the room is only 0.41%. From this and similar tests, patterns with 300 × 560 grid points for the room (300 in the vertical direction) and 60 × 160 for the chimney (60 × 320 for the long chimney of series b) are used. Also, post-solution checks of values of $\gamma^*$, the non-dimensional distance from closest wall, of grid points closest to the walls show $\gamma^*$ being well below 1, thus confirming that the grid patterns used are sufficiently fine for the Chien turbulence model.

Results and Discussion

Air-mass-flow rate and flow pattern have been considered for a range of configuration of the flow domain, and temperature of the chimney’s heated plate. The configuration’s factors include length and inclination of the chimney, its location and locations of the room’s other openings, and the presence of a vertical
partition in the room. The considered configurations (series) are shown in Table 1. Location of the chimney refers to location of the opening slot on the room's roof, which is also the lower entrance to the chimney. In all series, SC's length is 1.56 m = (3 - 0.3)/(2cos30°), except for series b* wherein this length is twice as long, at 3.12 m.

Figure 2 shows the net air-mass-flow rate \( m \) in terms of the temperature difference \( \Delta T \) between the chimney's heated plate and ambient air, for different configurations (series). On log-log scales, the \( m \)-versus-\( \Delta T \) curves are nearly straight lines, but with different slopes. The highest flow rates occur when chimney is located at a roof corner and inlet to the room at the opposite roof corner (series d, e), but only at sufficiently high \( \Delta T \) (in contrast to low-\( \Delta T \) results indicated by \( y \)-series). With inlets to the room at its lower corners, there is little difference in \( m \) (series t, h, u and z); 2 inlets (series l) give slightly more flow than a single inlet (series u); similarly with series l versus g (shown in figure 3). That maximum flow rates occur with series d and e is believed to be due to the longer and not-so-twisted flow path (as in series h, z) from room's inlets to the chimney, resulting in stronger stack effect (flow pattern is shown in figure 4).

Figure 2 also shows that a partition reduces \( m \) slightly (series e versus d). This agrees with expectation, as partitions offer extra resistance to the fluid motion.

Figure 3 shows an increasing \( m \) with configuration for \( \Delta T = 20^\circ C \). Values from series z versus h (figure 2) and z versus t indicate that a slightly higher \( m \) is obtained when chimney's upper plate is heated. But when chimney is upright, a left heated-plate (on the room side, series l) gives a much higher flow rate than the right one (series m), whose flow pattern (figure 4) shows much more back-flow at chimney exit than series l (not shown)). Series i gives maximum flow rate, much larger than that from series d and e. This is believed to be due to a taller vertical height rather than a larger internal passage, however. It's interesting to note that air flow in series I would be sufficient to supply an ample rate of 8 l/s to 19 persons in a room of 3-m length.

Figure 4 shows the net air-mass-flow rate \( m \) in terms of the temperature difference \( \Delta T \) between the solar chimney (SC)'s heated plate and ambient air.

Readings from series e and b* show that doubling the chimney's length results in near double in flow rate; thus \( m \) is nearly proportional to chimney's length.

Flow patterns corresponding to a number of series are shown in figure 4, with \( \Delta T = 20^\circ C \). Note that in this figure, the horizontal and 2 vertical lines inside the room should be disregarded (except for the section of the vertical line corresponding to the partition at A (see figure 1) when this is present); these lines are used in the discretisation of the computational domain, and thus are not part of the solution contours. The figure shows that the presence of a partition (series d) results in the lower part of the room being ventilated better. This is often a very desirable aspect in ventilation. This and other similar figures (not shown) also indicate that when there is significant back-flow at the chimney's exit, flow rate is correspondingly low; an example is series m (versus series I).
Figure 4. Flow pattern of some representative series, with $\Delta T = 20^\circ$C; (i) series m; (ii) series d; (iii) series i; (iv) series z.
Welcome

The Organising Committee of the 17th Australasian Fluid Mechanics Conference would like to take the opportunity to welcome you to New Zealand, to Auckland and to the 17th meeting of this conference.

The Australasian Fluid Mechanics Conferences have been held triennially since the series inception as the Australasian Hydraulics and Fluid Mechanics Conference at the University of Western Australia 1962. This Conference is the 17th meeting, the 4th of the series to be held in New Zealand, and the 3rd to be held at The University of Auckland. These conferences, at which all contributions are presented orally, have the tradition of providing graduate students with an opportunity to interact with a community of highly experienced researchers, and have played a very significant role in nurturing fluid mechanics research in Australasia and the surrounding region for nearly four decades.

The role of developing this vibrant community has been formalised with the recent formation of the Australian Fluid Mechanics Society, and this is the first event which has been associated with the AFMS. A highlight of this meeting will be the celebration of the induction of the 10 inaugural AFMS Fellows at the conference banquet.

A feature of the AFMC series is the broad range of fluid mechanics research presented, and this conference is no different as evidenced by the 24 themes identified in the programme, the 235 presentations and over 200 papers submitted. The flavour of a particular meeting is set by the plenary speakers and the conference committee is very grateful for the contributions of David Boger, Vladimir Nikora, Mike O'Sullivan, Andrew Pollard and Phil Schwarz for providing high quality thematic foc.

As Chair I would like to acknowledge here the efforts of John Carter the conference secretary and the other members of the conference Organising Committee: Rosalind Archer, John Chen, Richard Clarke, Stephen Coleman, Richard Flay, Bruce Malville, Roger Nokes, Stuart Norris, Peter Richards and Rajnish Sharma. I also acknowledge the 120 reviewers of the abstracts and papers. The Committee would particularly like to thank Tessa Hegemann, of the Conference Management unit in The University of Auckland’s Centre for Continuing Education who has tirelessly provided the administrative support that has allowed the committee to concentrate on the technical aspects of the conference.

Finally, the Committee would also like to acknowledge the support of the Faculty of Engineering at The University of Auckland and the Conference sponsors: the platinum sponsor Olympus, COMSOL Multiphysics, LasTEK & LaVision, LEAP Australia, Kenelec Scientific, Duff and Macintosh, New Spec & Spectra-Physics.

Gordon Mallinson
Chair
**17AFMC Programme Timetable**

**Monday**

**09:00 - 09:30**
- Welcome and The National Levels of Sustained Research
- Paper 278 Presentation

**09:30 - 10:30**
- Paper 278 - Presenting Author: Dr. Andrew Ooi
- Paper 279 - Presenting Author: Mr. Rajinesh Singh
- Paper 280 - Presenting Author: Ms. Sarah Preissler
- Paper 281 - Presenting Author: Dr. Biju Sahoo
- Paper 282 - Presenting Author: Mr. Cheng-Chang Chen
- Paper 283 - Presenting Author: Mr. Ming Zhou

**10:30 - 11:00**
- Coffee Break

**11:00 - 12:00**
- Paper 284 - Presenting Author: Dr. Andrew Ooi
- Paper 285 - Presenting Author: Mr. Rajinesh Singh
- Paper 286 - Presenting Author: Ms. Sarah Preissler
- Paper 287 - Presenting Author: Dr. Biju Sahoo
- Paper 288 - Presenting Author: Mr. Cheng-Chang Chen
- Paper 289 - Presenting Author: Mr. Ming Zhou

**12:00 - 13:00**
- Lunch Break

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<thead>
<tr>
<th>Paper</th>
<th>Title</th>
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<tr>
<td>278</td>
<td>Hybrid volume-of-fluid and dilute particle solver for oil-in-water emulsions</td>
<td>Dr. Andrew Ooi</td>
</tr>
<tr>
<td>279</td>
<td>Model for Droplet Spreading and Evaporation on Semi-confined Geometries</td>
<td>Mr. Rajinesh Singh</td>
</tr>
<tr>
<td>280</td>
<td>Development of a CFD model for an oscillating hydrofoil</td>
<td>Ms. Sarah Preissler</td>
</tr>
<tr>
<td>281</td>
<td>Microfluidic analysis of secondary vortices with observations of primary vortices in single bubble cavitation flutter and shock expansion</td>
<td>Dr. Biju Sahoo</td>
</tr>
<tr>
<td>282</td>
<td>A Two-Sided Model for Droplet Spreading and Evaporation on Semi-confined Geometries</td>
<td>Mr. Cheng-Chang Chen</td>
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<tr>
<td>283</td>
<td>Flow control in microchannels using grooves in confining walls</td>
<td>Mr. Ming Zhou</td>
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**13:00 - 14:00**
- Paper 284 - Presenting Author: Dr. Andrew Ooi
- Paper 285 - Presenting Author: Mr. Rajinesh Singh
- Paper 286 - Presenting Author: Ms. Sarah Preissler
- Paper 287 - Presenting Author: Dr. Biju Sahoo
- Paper 288 - Presenting Author: Mr. Cheng-Chang Chen
- Paper 289 - Presenting Author: Mr. Ming Zhou

**14:00 - 15:00**
- Paper 284 - Presenting Author: Dr. Andrew Ooi
- Paper 285 - Presenting Author: Mr. Rajinesh Singh
- Paper 286 - Presenting Author: Ms. Sarah Preissler
- Paper 287 - Presenting Author: Dr. Biju Sahoo
- Paper 288 - Presenting Author: Mr. Cheng-Chang Chen
- Paper 289 - Presenting Author: Mr. Ming Zhou

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**15:00 - 16:00**
- Paper 284 - Presenting Author: Dr. Andrew Ooi
- Paper 285 - Presenting Author: Mr. Rajinesh Singh
- Paper 286 - Presenting Author: Ms. Sarah Preissler
- Paper 287 - Presenting Author: Dr. Biju Sahoo
- Paper 288 - Presenting Author: Mr. Cheng-Chang Chen
- Paper 289 - Presenting Author: Mr. Ming Zhou

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**16:00 - 17:00**
- Paper 284 - Presenting Author: Dr. Andrew Ooi
- Paper 285 - Presenting Author: Mr. Rajinesh Singh
- Paper 286 - Presenting Author: Ms. Sarah Preissler
- Paper 287 - Presenting Author: Dr. Biju Sahoo
- Paper 288 - Presenting Author: Mr. Cheng-Chang Chen
- Paper 289 - Presenting Author: Mr. Ming Zhou

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**17:00 - 18:00**
- Paper 284 - Presenting Author: Dr. Andrew Ooi
- Paper 285 - Presenting Author: Mr. Rajinesh Singh
- Paper 286 - Presenting Author: Ms. Sarah Preissler
- Paper 287 - Presenting Author: Dr. Biju Sahoo
- Paper 288 - Presenting Author: Mr. Cheng-Chang Chen
- Paper 289 - Presenting Author: Mr. Ming Zhou

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**18:00 - 19:00**
- Paper 284 - Presenting Author: Dr. Andrew Ooi
- Paper 285 - Presenting Author: Mr. Rajinesh Singh
- Paper 286 - Presenting Author: Ms. Sarah Preissler
- Paper 287 - Presenting Author: Dr. Biju Sahoo
- Paper 288 - Presenting Author: Mr. Cheng-Chang Chen
- Paper 289 - Presenting Author: Mr. Ming Zhou

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**19:00 - 20:00**
- Paper 284 - Presenting Author: Dr. Andrew Ooi
- Paper 285 - Presenting Author: Mr. Rajinesh Singh
- Paper 286 - Presenting Author: Ms. Sarah Preissler
- Paper 287 - Presenting Author: Dr. Biju Sahoo
- Paper 288 - Presenting Author: Mr. Cheng-Chang Chen
- Paper 289 - Presenting Author: Mr. Ming Zhou

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**20:00 - 21:00**
- Paper 284 - Presenting Author: Dr. Andrew Ooi
- Paper 285 - Presenting Author: Mr. Rajinesh Singh
- Paper 286 - Presenting Author: Ms. Sarah Preissler
- Paper 287 - Presenting Author: Dr. Biju Sahoo
- Paper 288 - Presenting Author: Mr. Cheng-Chang Chen
- Paper 289 - Presenting Author: Mr. Ming Zhou

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**21:00 - 22:00**
- Paper 284 - Presenting Author: Dr. Andrew Ooi
- Paper 285 - Presenting Author: Mr. Rajinesh Singh
- Paper 286 - Presenting Author: Ms. Sarah Preissler
- Paper 287 - Presenting Author: Dr. Biju Sahoo
- Paper 288 - Presenting Author: Mr. Cheng-Chang Chen
- Paper 289 - Presenting Author: Mr. Ming Zhou

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**22:00 - 23:00**
- Paper 284 - Presenting Author: Dr. Andrew Ooi
- Paper 285 - Presenting Author: Mr. Rajinesh Singh
- Paper 286 - Presenting Author: Ms. Sarah Preissler
- Paper 287 - Presenting Author: Dr. Biju Sahoo
- Paper 288 - Presenting Author: Mr. Cheng-Chang Chen
- Paper 289 - Presenting Author: Mr. Ming Zhou