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# CuO single crystal with exposed {001} facets - A highly efficient material for gas sensing and Li-ion battery applications

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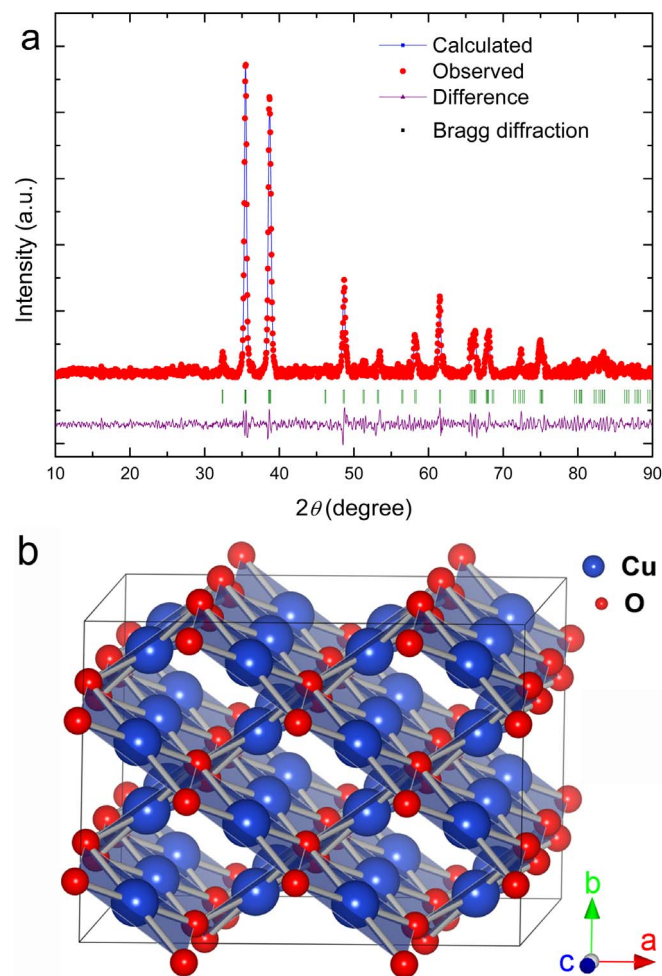
Single crystal copper oxide nanoplatelets with a high percentage of {001} facets were synthesized by a facile hydrothermal approach. The as-prepared materials were characterized by X-ray diffraction, Fourier transform infrared spectroscopy, field emission scanning electron microscopy, and high resolution transmission microscopy. Via density functional theory calculations, it was found that the {001} facets are active crystal planes. When the single crystal CuO nanoplatelets were applied as an anode material in Li-ion batteries, they demonstrated outstanding electrochemical performance with high lithium storage capacity, satisfactory cyclability, and excellent high rate capacity. When used as a sensing material in gas sensors, they exhibited a superior sensitivity towards toxic and flammable gases.

Nanosize materials have peculiar properties that are not expected in the bulk phase, and elucidation of these properties has already led to breakthroughs in various fields of science and technology. The electrical and optical properties of nanoparticles are often size- and shape-dependent, and hence, it is essential to be able to control the particle size, shape, and distribution<sup>1–3</sup>. These require a detailed understanding of the mechanisms of nucleation and growth, as well as processes such as aggregation and coarsening<sup>4,5</sup>. Compared to the morphogenesis of metal oxide, however, much less attention has been paid to the crystal plane effects of nanocrystals in different reactions. Actually, they should have attracted remarkable attention, because they are essential factors for the interaction<sup>6–9</sup>. Uniform crystal planes at the nanoscale are the most important prerequisite in investigating the differences in properties due to facets from different planes. As one example, TiO<sub>2</sub> with exposed {001} facets has been widely researched<sup>10–14</sup>, because both theoretical and experimental studies have found that the minority {001} facets in the equilibrium state are especially reactive<sup>15</sup>.

As an important transition metal oxide, cupric oxide (CuO) has been extensively studied due to its diverse applications, such as in gas sensors, catalysis, batteries, high critical temperature superconductors, solar energy conversion, and field emission emitters<sup>16–19</sup>. As anode for Li-ion battery, CuO can achieve superior electrochemical performance<sup>20</sup>. In this study, we found that the {001} facets of CuO are more reactive. We show that *p*-type semiconductor gas sensors based on CuO with {001} exposed facets have a rapid response and high sensitivity to very low concentrations of flammable and toxic gases. Furthermore, CuO with {001} exposed facets has also demonstrated significantly improved lithium storage capacity and cycling stability in lithium ion cells. In particular, it has excellent high rate performance in Li-ion batteries.

## Results and discussion

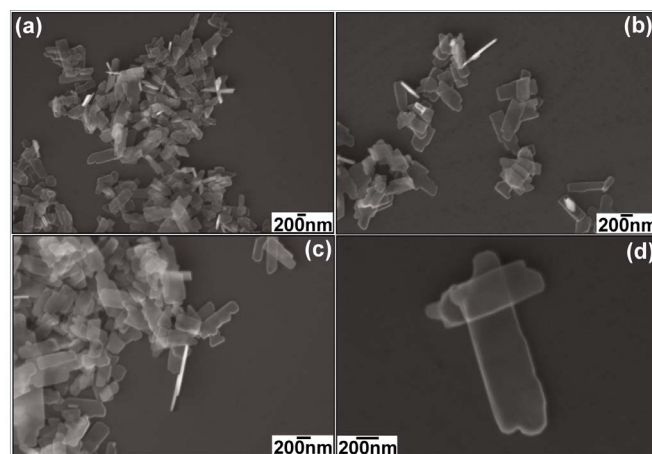
**Structural and morphological analyses.** Figure 1(a) shows the refined X-ray diffraction pattern of the CuO nanoplatelets. It can be well indexed with a pure monoclinic symmetry unit cell with space group C2/c and has satisfactory convergence factors ( $R_{wp} = 9.23\%$ ,  $R_p = 7.57\%$ ,  $\chi^2 = 3.17$ ). The lattice parameters were refined as listed in Table 1. The refined crystal structure of CuO is shown in Figure 1(b), in which the Cu atom is coordinated by four coplanar O atoms forming an almost rectangular parallelogram, while the O coordination polyhedron has four Cu atoms at the corners of a distorted tetrahedron. The CuO, by sharing opposite edges, forms two ribbons of parallelograms running in the [110] and  $[\bar{1}\bar{1}0]$  directions. In the [101] and  $[\bar{1}0\bar{1}]$  directions, two types of -Cu-O-Cu- chains could be differentiated.



**Figure 1** | (a) Rietveld refinement pattern of X-ray diffraction data for as-prepared CuO nanoplatelets. (b) Crystal structure of the refined CuO unit cell.

Upon energy minimisation of the bulk crystal, the lattice parameters of CuO were calculated, as shown in Table 1. All the differences between the calculated results and the experimental values are less than 1%, indicating good accuracy of the computational method.

The morphology of as-prepared CuO is revealed by FESEM, as shown in Figure 2. In the low magnification FESEM image (Figure 2(a)), it can be seen that the CuO consists of well-defined nanoplatelet structures with irregular, roughly rectangular outlines. The width of the cross-section is between 200 and 300 nm. The thickness is less than 50 nm, as shown in Figure 2(b) and (c). Because the nanoplatelets are thin enough, they are transparent even when they overlap, which is demonstrated by the high magnification FESEM image (Figure 2(d)). For comparison, different alkali and amount of the alkali were used for preparation. When the NaOH solution was used instead of the LiOH, CuO nanoparticles were



**Figure 2** | FESEM images of the as-prepared faceted CuO nanoplatelets: (a) low magnification, (b) and (c) medium magnification, and (d) high magnification.

obtained as shown by Figures S1 and S2 in the Supporting Information (SI). The morphology of as-prepared CuO nanoparticles is not changed too much if the amount of NaOH was reduced, as shown by Figure S3 (a) (SI). When the LiOH amount was increased to 20 mmol, it was found that the as-prepared CuO still preserved the nanoplatelets morphology as shown by Figure S3(b) (SI). Therefore, the kind of alkalis is more important for the synthesis of CuO nanoplatelets.

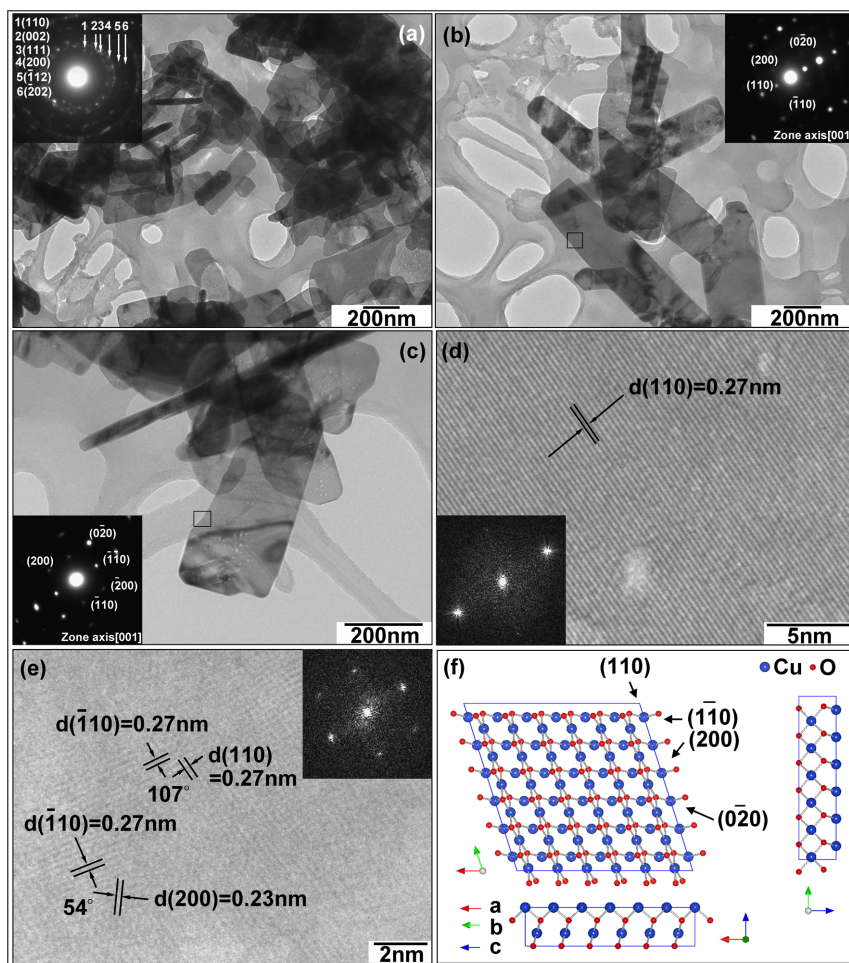
The morphology of the as-prepared CuO nanoplatelets was further characterized by TEM and HRTEM analyses (Figure 3). Figure 3(a) contains a low magnification TEM image of the CuO nanocrystals and its corresponding SAED pattern (inset in Figure 3(a)). From the SAED pattern, it can be seen that all diffraction rings can be indexed to the monoclinic CuO phase. High magnification TEM images (Figure 3(b) and (c)) demonstrate that the free-standing nanocrystals are thin enough to be transparent. The SAED patterns (insets in Figure 3(b) and (c)) taken from different CuO nanocrystals show the same rhombus of diffraction spots along the [001] zone axis (as the (200), (020), (011), and (110) crystal planes can be indexed), suggesting that each free-standing CuO nanocrystal has the same {001} facets exposed. The lattice resolved HRTEM images (Figures 3(d) and (e)) directly present the (110), ( $\bar{1}10$ ), and (200) crystal planes with 0.27 nm and 0.23 nm  $d$ -spacings, and  $107^\circ$  and  $54^\circ$  interfacial angles. The normal direction to these crystal planes can be further confirmed as [001]. Thus, as illustrated by Figure 3(f), the as-prepared CuO faceted single-crystal has {001} dominant exposed crystal planes.

The quality of the CuO nanoplatelets was also examined by FTIR spectroscopy, as shown in Figure S4 (SI). The typical FTIR spectrum of the as-prepared CuO nanoplatelets has sharp and narrow peaks at 663, 574, and 458  $\text{cm}^{-1}$ , characteristic peaks for the Cu-O stretching vibration, which is redshifted compared with a previous report<sup>20</sup>. The high-frequency mode at 663  $\text{cm}^{-1}$  is the Cu-O stretching vibration along the [101] direction, and the mode at 574  $\text{cm}^{-1}$  is the Cu-O stretching vibration along the [011] direction<sup>21</sup>.

To study the surface activities of CuO, a set of suitable surface energies for the crystal planes were calculated. The surface region in our calculations is composed of a finite number of two-dimensional infinite planes formed by cutting the crystal along a particular Miller index (hkl) plane. In each plane, a two-dimensional cell represents every site in the plane. Following the approach of Tasker<sup>22</sup>, several of these cells in successive planes comprise the basic repeated unit of CuO that has the composition of the bulk crystal unit cell. We treated the {100}, {010}, {001}, {110}, {101}, and {011} crystal planes in this work. According to the previous results<sup>23,24</sup>, the relaxed surface

**Table 1** | Optimized lattice parameters of CuO

Lattice parameters	CuO		
	Calculated	Experimental	Difference/%
$a/\text{\AA}$	4.68	4.69	0.21
$b/\text{\AA}$	3.42	3.43	0.29
$c/\text{\AA}$	5.14	5.10	0.78
$\alpha, \beta, \gamma/^\circ$	90 99.53 90	90 99.51 90	0.02
Cell volume/ $\text{\AA}^3$	80.59	81.37	0.96



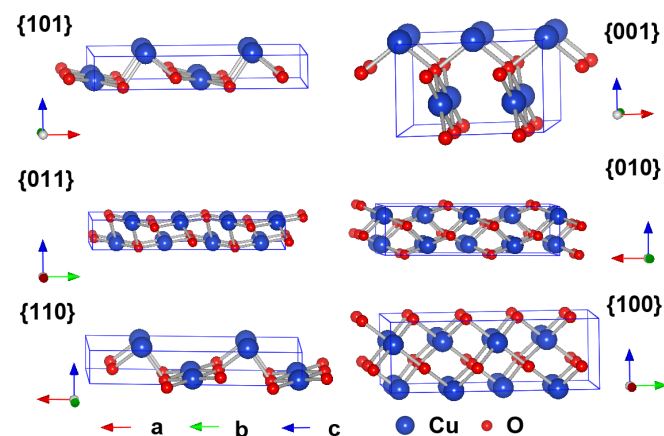
**Figure 3** | (a) Low magnification TEM image of CuO. (b–c) High magnification TEM images of CuO. (d–e) Lattice resolved HRTEM images of CuO. (f) Crystal structure of {001} crystal planes of CuO. Insets in (a), (b), and (c) are selected area electron diffraction (SAED) patterns. Insets in (d) and (e) are fast-Fourier-transform (FFT) patterns.

energy is lower than the unrelaxed surface energy, which indicates the importance of modelling the surface relaxation. Herein, the ion coordinates in each cell were calculated by using the relaxed crystal coordinates. The relaxed unit cell of CuO was extended to construct the surface models as shown in Figure 4.

The most stable surface energies for the relaxed surface regions of CuO are listed in Table 2. It shows that the {010} surfaces of CuO have the lowest energy ( $1.04 \text{ J m}^{-2}$ ), while the {011} surfaces have the second smallest surface energy, followed by the {100} surfaces, which are all less than  $2 \text{ J m}^{-2}$ , indicating their structural similarities. {101} and {110} surfaces have relatively larger surface energies, which are all more than  $2.5 \text{ J m}^{-2}$ . {001} surfaces have the largest surface energy,  $4.23 \text{ J m}^{-2}$ . Therefore, the most active surfaces of CuO are {001}. We also calculated the interaction energy between lithium and different facets of CuO. It can be seen that {001} facets yield the lowest interaction energy ( $1.88 \text{ J m}^{-2}$ ), while {010} facets show the highest value ( $8.47 \text{ J m}^{-2}$ ). Apparently, the {001} crystal planes show the highest activity due to their high surface energy, which can provide more reactive sites to facilitate interaction with Li, and therefore, they consume the lowest energy to accomplish reactions. We expect that the as-prepared CuO nanoplatelets with exposed high energy {001} facets could provide highly reactive sites, leading to enhanced performance for nanotechnological applications. We tested the gas sensing and Li-ion battery properties of the CuO nanoplatelets.

**Gas-sensing performance.** Figure 5 shows the gas sensitivities of CuO nanoplatelet based sensors as a function of the gas

concentrations of acetone, butanol, ethanol, and isopropanol. For comparison, the sensitivities of CuO nanoparticles are also given in Figure 5. Generally, the sensitivity of the CuO nanoplatelets is higher than that of the CuO nanoparticles. The CuO nanoplatelets presented outstanding sensitivity to low concentrations of gases. The insets in Figure 5 show the real-time responses towards different gases. Ten testing cycles were recorded, corresponding to



**Figure 4** | Side views of relaxed {101}, {001}, {011}, {010}, {110}, and {100} crystal planes of CuO.





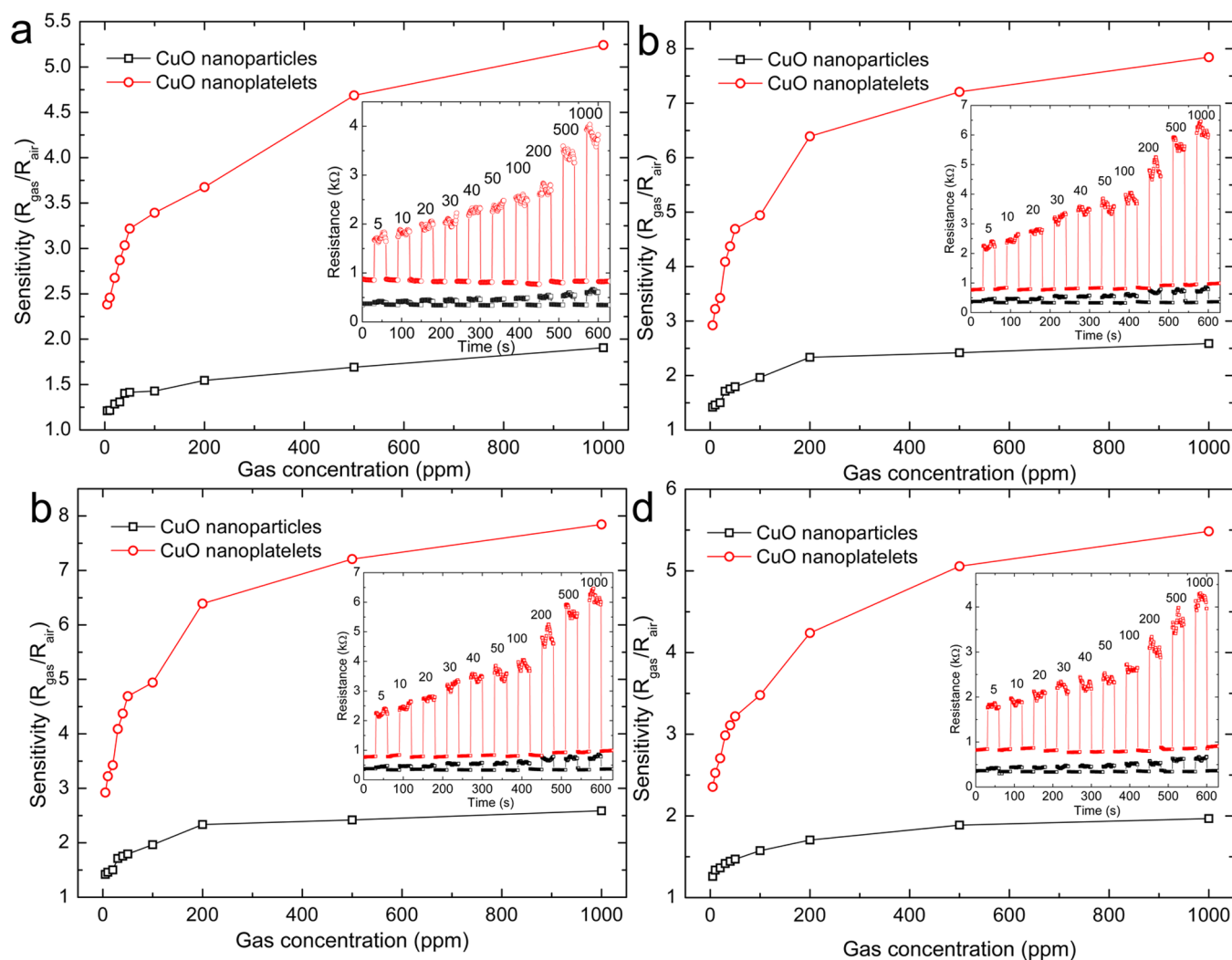
**Table 2 |** Calculated surfaces energies, interaction energy between lithium and different facets of CuO

Surface	CuO		
	$d_{\text{spacing}}, d_{\text{hkl}}/\text{\AA}$	Surface energy/ $\text{J m}^{-2}$	Interaction energy/ $\text{J m}^{-2}$
001	5.038	4.23	1.88
010	3.41	1.04	8.47
100	4.589	1.73	6.05
011	2.824	1.22	5.42
101	3.144	2.51	3.42
110	2.737	2.69	3.43

ten different concentrations from 5 to 1000 ppm. It can be seen that the electric resistance of the sensor ( $R_{\text{sensor}}$ ) increased abruptly on the injection of gas, and then decreased rapidly and recovered to its initial value after the test gas was released. The response magnitude of the CuO nanoplatelets increased significantly with increasing concentration of the test gas and was much higher than that of the CuO nanoparticles. After many cycles, the resistance of the sensor recovered to its initial state, which indicates that the CuO

nano-platelet sensor has good reversibility. The response time and recovery time (defined as the times required to reach 90% of the final equilibrium value) of the CuO nanoplatelet sensor were both less than 1 s. It was also found that the on and off responses could be repeated after continuous measurement for two weeks, with no detectable changes in the signal, illustrating the good reversibility and stability of the CuO nanoplatelet based sensor.

From Ohm's law, the electrical resistance of the sensor underwent a decreasing and increasing process when the test gas was turned on and off, respectively. The mechanism responsible for the gas sensing is considered to follow the surface conduction modulation<sup>25,26</sup>. CuO is a *p*-type semiconductor, in which the positive holes are the majority carriers<sup>43</sup>. Upon exposure to a reducing gas, the density of *p*-type charge carriers (positive holes) would decrease due to surface adsorption and chemical reactions between the gas and the oxygen adsorbates (electron acceptors), resulting in an increase of the sensor's resistance<sup>27</sup>. The improved sensing performance of the CuO nanoplatelets should be ascribed to the high surface energy of the predominantly exposed {001} crystal planes, which can provide more active sites and improve the kinetics of both the reactions of the reducing gases with surface-adsorbed oxygen and the replacement of the latter from the gas phase.



**Figure 5 |** Gas sensitivities of CuO nanoplatelet based sensors as a function of the gas concentrations of (a) acetone, (b) butanol, (c) ethanol, and (d) isopropanol gases. For comparison, the gas sensing performances of CuO nanoparticles are also presented. The insets are the corresponding real-time gas-sensing curves for the CuO nanoplatelets (red) and nanoparticles (black), with the concentrations of the test gas in ppm given above the peaks.



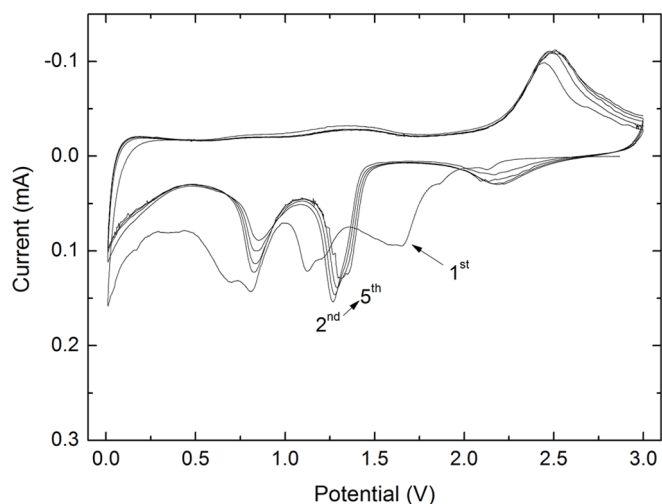
**Electrochemical performance of CuO nanoplatelets.** The electrochemical performance of the CuO nanoplatelets was investigated by cyclic voltammetry (CV) and galvanostatic charge and discharge measurements. Figure 6 shows the CV curves. They are similar to previous reported results<sup>28,29</sup>. During the first cycle, three cathodic peaks are observed at 1.6, 1.2, and 0.8 V, respectively. These peaks can be associated with a multistep reaction related to the creation of a  $\text{Cu}_{1-x}\text{Cu}_{1-x/2}\text{O}_{1-x/2}$  ( $0 \leq x \leq 0.4$ ) solid solution within the CuO phase, the formation of a  $\text{Cu}_2\text{O}$  phase, and finally, the decomposition into metallic Cu and  $\text{Li}_2\text{O}$ , respectively<sup>30,31</sup>. Furthermore, a relatively strong cathodic peak at low potential between 0.01 and 0.3 V can be derived from the formation of a solid electrolyte interphase (SEI), which can lead to an extra capacity besides the reversible Li driven decomposition of the transition-metal oxide<sup>32</sup>. The anodic peak appearing at 2.5 V is due to the oxidation of metallic Cu-Cu (I) and to the partial oxidation of Cu (I)-Cu (II). The broad bump around 1.45 V could be associated with decomposition of the SEI layer<sup>30,31</sup>. In the subsequent cycles, the cathodic and anodic peaks have a positive shift. Both the peak current and the integrated area of the anodic peak are stable and well overlapped with each other, indicating the high reversibility and good capacity retention of the CuO nanoplatelets. For comparison, the CV curves of the CuO nanoparticles are given in Figure S5 (SI). There is a broad peak that appears at about 1.03 V in the cathodic process in the first cycle, which could be attributed to the reduction process from Cu (II) to Cu. In addition, this cathodic process is also associated with electrolyte decomposition and the reversible conversion reaction of lithium ion intercalation to form  $\text{Li}_2\text{O}$ . An anodic peak is present at about 2.42 V, corresponding to the reversible oxidation of Cu to Cu (II). During the anodic process, both the peak current and the integrated area of the anodic peak are decreased, indicating capacity loss during the charging process.

The CV peaks and their variations are consistent with the galvanostatic discharge-charge profiles as shown in Figure 7. Three distinct potential slopes are visible in both the discharge and the charge curves, as shown in Figure 7(a). As indicated by the arrows in Figure 7(a), in the discharge process, there are three different plateaus, located at 2.16, 1.5, and 0.85 V, respectively, while in the charge process, the slopes are located at around 1.26, 2.2, and 2.56 V. In the first cycle at a current density of  $134 \text{ mA g}^{-1}$  (0.2 C), the CuO nanoplatelet electrode achieved a discharge capacity of  $1323 \text{ mA h g}^{-1}$ . The subsequent discharge profiles show a more pronounced step-like profile, with a typical upwards shift of the

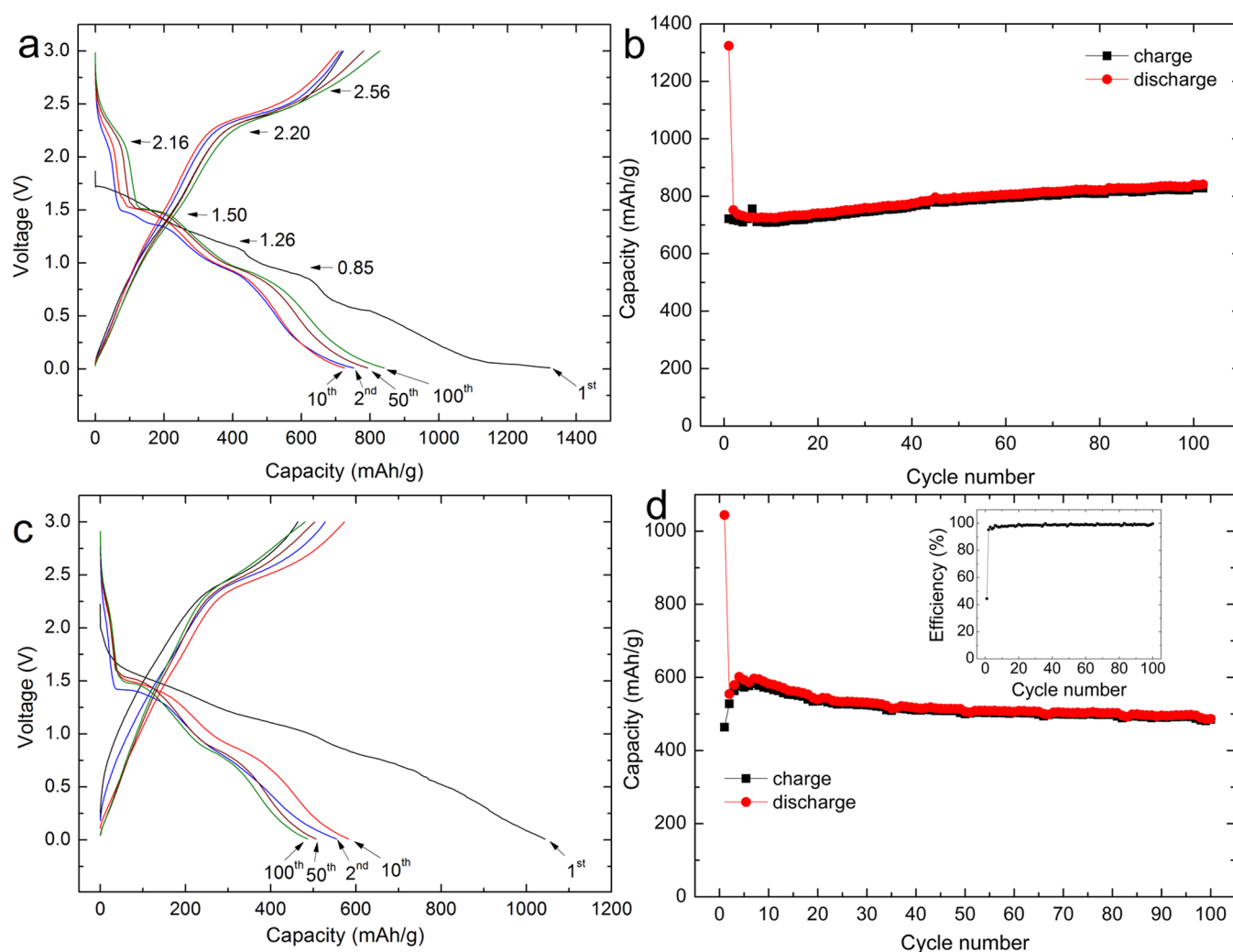
voltage plateau due to the enhanced tendency towards dissolution of the regenerated nano-CuO after the initial reaction with Li, as discovered for the CuO/Li primary battery<sup>33</sup>. Figure 7(b) presents the evolution of the charge and discharge capacities during cycling. It can be seen that after 100 cycles, the test cell delivered a discharge capacity of  $841 \text{ mA h g}^{-1}$ , demonstrating ultra-high cyclability at a low current rate. For comparison, the initial discharge capacity of the CuO nanoparticle electrode was  $1050 \text{ mA h g}^{-1}$ , and from the second cycle, the discharge capacity decreased continuously and only delivered  $340 \text{ mA h g}^{-1}$  after 100 cycles (Figure S6, SI). Obviously, the CuO nanoplatelet electrode shows better capacity retention. On increasing the current rate to 0.5 C, the capacity is still as high as  $700 \text{ mA h g}^{-1}$  (Figure S7, SI). The extra capacity above the theoretical value ( $670 \text{ mA h g}^{-1}$ )<sup>29</sup> is most likely due to the decomposition of electrolyte during the discharge process. In this step, the SEI layer is formed and covers the CuO to protect the electrode from the intercalation of solvent<sup>34,35</sup>. Figure 7(c) and (d) shows the electrochemical performance of the CuO nanoplatelet electrode at the 1 C current rate. It can be seen that the initial discharge capacity is  $1041 \text{ mA h g}^{-1}$ . Although the discharge capacity is reduced to  $552 \text{ mA h g}^{-1}$  in the second cycle, from the third cycle, the coulombic efficiency reaches as high as 98.96% (inset in Figure 7(d)), and the electrode presents remarkable cyclability ( $502 \text{ mA h g}^{-1}$  after 100 cycles).

We also cycled CuO nanoplatelets at the high current rates of 5 C and 10 C for 1000 cycles to investigate their realistic applicability. As shown in Figure 8(a), the charge and discharge capacity curves are almost overlapping due to the ultra-high coulombic efficiency, with an average value of 99.27% over 1000 cycles. After 200 cycles at 5C, there is rapid capacity increase. Such capacity increase phenomenon was also reported by other publications on CuO<sup>20,36,37</sup>, other metal<sup>38</sup>, and metal oxide anodes<sup>39,40</sup>. It should be attributed to the reversible formation of gel-like polymeric species in the SEI films due to the catalytic activity of metals in the anodes<sup>20,36–42</sup>, and the volume change of the CuO anodes during the lithium-ion insertion/extraction may build more reversible polymeric species, which result into the continuous increase in capacity after 200 cycles<sup>20</sup>. It might be also the reason that the capacity of CuO nanoplatelets at 5C is a little lower than that of 10C rate in the first 100 cycles, because more reversible polymeric species were generated at 10 C rate in the first 100 cycles, which result into the higher capacity. The cells retained above 751 and 309  $\text{mA h g}^{-1}$  discharge capacities at 5 C and 10 C, respectively, after 1000 cycles. CuO nanoparticle electrodes were also tested at 5 C and 10 C. Although they presented high initial discharge capacity, they just retained 150 and 68  $\text{mA h g}^{-1}$  discharge capacity after 1000 cycles, respectively (Figure S8, SI). It can be seen that the as-prepared CuO nanoplatelet electrode exhibited superior electrochemical performance in long-term cycling at high current rate. These results are the better than other reported CuO anode materials with different size<sup>36,44,45</sup>. Figure 8(b) compares the cycling performances of CuO nanoplatelets and CuO nanoparticles at varied current rates. The CuO nanoplatelets yield 914, 798, 701, 601, 516, 411, and 322  $\text{mA h g}^{-1}$  discharge capacity at the 0.1, 0.2, 0.5, 1, 2, 5, and 10 C current rates, respectively, and when the current rate is reversed back to the initial low current rate, the cell capacity recovers to the original value, indicating that the integrity of the CuO nanoplatelets has been preserved, even after high rate cycling. This implies that the CuO nanoplatelets are tolerant to varying charge and discharge currents, which is preferred for high power applications. In contrast, when the CuO nanoparticle electrode was operated at elevated current rates, the capacity was decreased to an extremely low level. (At 10 C, the capacity is only 46  $\text{mA h g}^{-1}$ .)

The superior electrochemical performance of as-prepared CuO nanoplatelets should be ascribed to the predominantly exposed {001} facets. Generally, high energy surfaces have a large density of low-coordinated atoms situated on steps and kinks, giving them high reactivity<sup>46</sup>. This favours fast ion transfer between the surface and the



**Figure 6** | CV curves for the first 5 cycles of the CuO nanoplatelets in a lithium ion cell.



**Figure 7** | Discharge and charge profiles and cycling performance of CuO nanoplatelet electrode at 0.2 C (a, b) and 1 C (c, d) current rates. The inset in (d) is the corresponding coulombic efficiency.

interior<sup>47,48</sup>. Because the {001} facets have relatively high surface energy, they can provide reactive sites for the reaction with lithium ions, which can facilitate a fast conversion reaction towards lithium during the charge and discharge processes. On the other hand, the thin nanoplatelet architecture provides a short path for lithium transport and high surface area for interfacial lithium storage<sup>35,49</sup>.

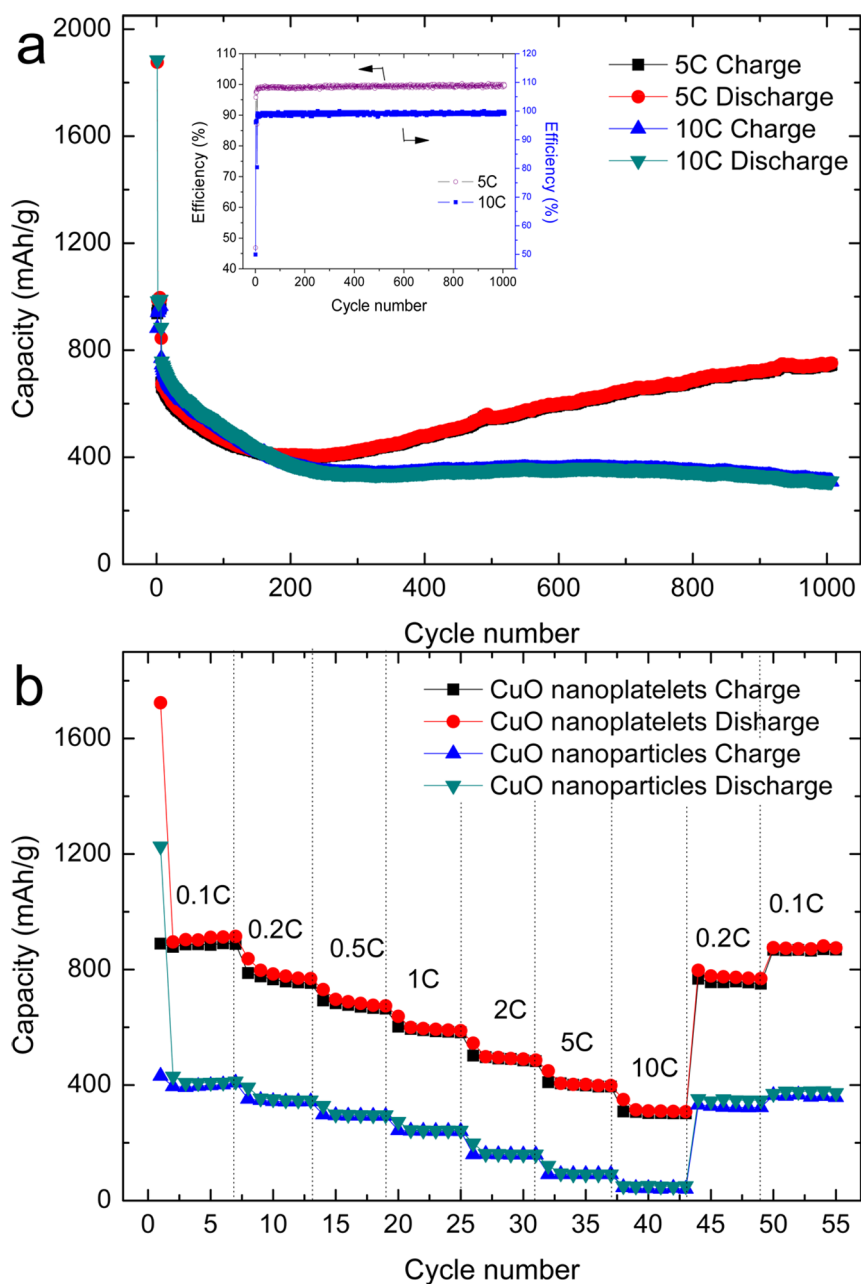
## Conclusion

CuO nanoplatelets were prepared by a simple low temperature hydrothermal method. XRD, FESEM, and HRTEM analyses revealed that the CuO nanoplatelets have predominantly exposed {001} crystal planes. Through the DFT calculations, it was found that the {001} facets have the highest surface energy. For this reason, the CuO nanoplatelets exhibited high sensing responses towards ethanol, acetone, butanol, and isopropanol, indicating their potential application in monitoring toxic and flammable gases. When used as anodes for lithium storage in Li-ion batteries, the CuO nanoplatelets delivered a higher reversible lithium storage capacity than the corresponding nanoparticles. After 100 cycles, the nanoplatelet electrode retained 841 and 502 mA h g<sup>-1</sup> capacity at the 0.2 C and 1 C rates, respectively. Furthermore, this material also demonstrated an outstanding high rate performance. The superior electrochemical performance should be ascribed to the predominantly exposed {001} reactive facets and the thin nanoplatelet architecture.

## Experimental Section

**Synthesis process.** Copper oxide nanoplatelets with {001} facets exposed were synthesized by a low temperature hydrothermal method. In a typical synthesis process, 1 mmol copper (II) chloride dihydrate (CuCl<sub>2</sub>·7H<sub>2</sub>O, 99+%, Sigma-Aldrich), and 3 mmol lithium hydroxide monohydrate (LiOH·H<sub>2</sub>O, Sigma-Aldrich) were first dissolved in 20 ml distilled water and stirred for ten minutes. The solution was then sealed in a Teflon-lined stainless steel autoclave (25 mL capacity). The autoclave was heated and maintained at 180°C for 2 h. The autoclave was then cooled down to room temperature, leading to the conversion of the precipitate into a black solid product. The precipitates were retrieved by washing with ethanol and water several times and drying in a vacuum oven at 70°C for 12 hours.

For comparison, CuO nanoparticles were synthesized: 1 mmol of CuCl<sub>2</sub>·2H<sub>2</sub>O was dissolved in 20 ml H<sub>2</sub>O, which was stirred with a magnetic stirrer. This solution was stirred for 15 min to ensure that the CuCl<sub>2</sub> was dissolved completely. Then, 10 ml 2 M NaOH aqueous solution was added dropwise into the CuCl<sub>2</sub> solution under constant stirring. The solution was then sealed in a Teflon-lined stainless steel autoclave, heated, and maintained at 200°C for 12 h. When the autoclave was cooled down to room temperature, the black CuO precipitate was washed with distilled water several times and then with ethanol, filtered, and dried in an oven at 70°C for 12 hours.



**Figure 8** | (a) Cycling performance of CuO nanoplatelet electrode at the high current rates of 5 C and 10 C. The inset shows the coulombic efficiency. (b) Rate performance of the CuO nanoplatelet and CuO nanoparticle electrodes.

**Structural and physical characterization.** The crystal structure and phase of the as-prepared materials were characterized by X-ray diffraction (XRD, Siemens D5000) using Cu K $\alpha$  radiation with  $2\theta$  ranging from  $20^\circ$  to  $90^\circ$  and a scanning step of  $0.02^\circ \text{ sec}^{-1}$ . The morphology was analyzed by high resolution field emission scanning electron microscopy (FESEM, Zeiss Supra 55VP). The microscope was operated at a working distance of 2 mm with an acceleration voltage of 20 kV, and an in-lens detector was used for the imaging. The details of the structure were further characterized by transmission electron microscopy (TEM) and high-resolution transmission electron microscopy (HRTEM, JEOL JEM-2011). Selected area electron diffraction (SAED) patterns were recorded by a Gatan charge-coupled device (CCD) camera in a digital format. The Fourier transform infrared (FTIR) spectra of the samples were collected using a Bruker Tensor 27 IR spectrometer with KBr as dispersant.

**Gas-sensing measurements.** The gas-sensing properties of the as-prepared faceted CuO were measured by a computer-controlled WS-30A gas-sensing measurement system. Figure S9 (SI) contains a schematic diagram of the system. In the procedure for making the sensors, the as-prepared CuO is first dispersed in tetraethylammonium tetrafluoroborate (Sigma-Aldrich, 99%, as binder) and ethanol to form a slurry. The mixture is then deposited as a thin film on a ceramic tube with Au electrodes and Pt conducting wires. Finally, the ceramic tube is heated at  $400^\circ\text{C}$  for 2 h to evaporate the solvent. The gas sensing measurements were carried out at a working temperature of  $150^\circ\text{C}$  and 30% relative humidity (RH).

**Electrochemical testing.** The electrodes were prepared by dispersing the as-prepared CuO (80 wt.%), acetylene carbon black (10 wt.%), and poly(vinylidene fluoride) binder (PVDF, 10 wt.%) in N-methyl-





2-pyrrolidone (NMP) to form a slurry. The resultant slurry was pasted onto copper foil using a doctor blade and dried at 100 °C for 12 h under vacuum conditions, followed by pressing at 200 kg cm<sup>-2</sup>. Electrochemical measurements were carried out using two-electrode coin cells with lithium metal as the counter electrode. The CR2032-type coin cells were assembled in an argon-filled glove box (UniLab, Mbraun, Germany). The electrolyte solution was 1 M LiPF<sub>6</sub> dissolved in a mixture of ethylene carbonate (EC) and dimethyl carbonate (DMC) with a volume ratio of 1 : 1. Cyclic voltammetry (CV) was carried out on a CHI 660C electrochemistry workstation with a scan rate of 0.1 mV s<sup>-1</sup> from 0.01 to 3.0 V in a two-electrode system. The charge-discharge measurements were performed at ambient temperature at different current densities in the voltage range from 0.01 to 3.0 V.

**Computational methods.** Electronic structure calculations were based on the density functional theory (DFT) + U approach<sup>50</sup>. The exchange-correlation energy functional was represented by the local-density approximation (LDA)<sup>51</sup> employing ultra-soft pseudopotential (USPP) formalism<sup>52</sup>. The valence configurations of the pseudopotentials were 2s2p4 for O and 3d104s1 for Cu. We used an energy cut-off of 350 eV in the plane-wave basis set expansion. Monkhorst-Pack k-point sets of 6 × 6 × 6 were used for an 8-atom unit cell of monoclinic CuO (space group C2/c). The on-site Coulomb repulsion (Hubbard U) was applied for Cu d states. The effective value for U of 9.0 eV for Cu was derived by fitting to experimental oxidation enthalpies, and it gave better estimates for band gaps<sup>53</sup>. The maximum self-consistent field convergent tolerance was less than 2 × 10<sup>-6</sup> eV atom<sup>-1</sup>. All calculations were performed in reciprocal space.

The surface energy per unit area,  $E_{\text{surface}}^{\text{hkl}}$ , of a particular surface is calculated from the difference between the energy of the surface block,  $E_{\text{surface block}}$ , and the energy of the same number of bulk ions,  $E_{\text{bulk}}$ , per unit area, A, and thus

$$E_{\text{surface}}^{\text{hkl}} = [E_{\text{surface block}} - E_{\text{bulk}}] / A \quad (1)$$

The variations in energy of the Li interactions on the different facets of the CuO,  $E_{\text{Li-CuO reactions}}^{\text{hkl}}$ , is calculated from the difference between the (hkl) surface energy,  $E_{\text{surface}}^{\text{hkl}}$ , the energy of Li monomer in vacuum,  $E_{\text{Li}}$ , and the total energy of the interacting system of the (hkl) CuO surface substrate and Li, per unit area, and thus

$$E_{\text{Li-CuO reactions}}^{\text{hkl}} = [E_{\text{surface+Li}}^{\text{hkl}} - (E_{\text{surface}}^{\text{hkl}} + E_{\text{Li}})] / A \quad (2)$$

For model construction, the depths of the surface regions were chosen to be large enough to ensure full relaxation of the surface ions (approximately 200 ions) and convergence of the surface energy. In each case, surface structures were fully relaxed until the total energy difference had converged within 0.001 eV.

We considered the low index surfaces containing planes that have the largest interplanar spacing ( $d_{\text{hkl}}$ ) and, as a result, are often the most stable. These are the most important planes morphologically predicted by the Bravais-Friedel-Donnay-Harker theory<sup>54</sup>. This approach assumes that the morphologically dominant faces (slowest growing) are those with thicker growth slices. Clearly, this is an over simplification, but it is nonetheless useful in reducing an infinite number of possibilities to a small set which includes the facets that will appear in the morphology.

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## Author contributions

D.-W.S. performed the experiments. X.-Q.X. assisted the experiments. S.-X.D. and G.-X.W. conceived the study. D.-W.S. wrote the manuscript. All authors discussed the results and reviewed the manuscript.

## Additional information

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